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CONTINUITY OF WEAK SOLUTIONS TO CERTAIN SINGULAR PARABOLIC EQUATIONS

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ABSTRACT

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It is demonstrated that weak solutions of (1.1) in the introduction are continuous in their domain of definition. The continuity up to the boundary is also investigated.

AMS(MOS) Subject Classifications: 35Kl0, 35Kl5, 35K20, 35K65

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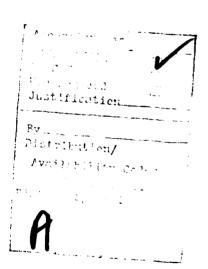
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SIGNIFICANCE AND EXPLANATION

The singular parabolic equations treated here serve as a model of heat conduction in processes where a change of phase occurs, such as water-ice, solidification of alloys, melting of metals.

Usually solutions of boundary value problems associated with these equations are found in a global sense, i.e. they are defined as equivalence classes in certain Sobolev spaces. It is of interest to decide whether they may be defined pointwise and if they possess some local regularity such as continuity.

In this paper we prove that global (weak) solutions are in fact continuous. Moreover we study under what circumstances the continuity can be extended up to the boundary of the domain where the process takes place.



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CONTINUITY OF WEAK SOLUTIONS TO CERTAIN SINGULAR PARABOLIC EQUATIONS

Emmanuele Di Benedetto

1. Introduction:

In this paper we study the continuity of weak solutions of parabolic "equations," with principal part in divergence form, of the type

(1.1)
$$\frac{\partial}{\partial t} \beta(u) - \operatorname{div} \dot{\tilde{a}}(x,t,u,\nabla_{x}u) + b(x,t,u,\nabla_{x}u) = 0$$

in the sense of distributions over a domain Q in \mathbb{R}^{N+1} . Here $\beta(\cdot)$ represents a maximal monotone graph in $\mathbb{R}\times\mathbb{R}$ such that $0\in\beta(0)$, $\frac{1}{2}$ is a map from \mathbb{R}^{2N+2} into \mathbb{R}^{N} and b maps \mathbb{R}^{2N+2} into \mathbb{R}^{1} .

Beside their intrinsic interest, inclusions such as (1.1) arise as a model to a variety of diffusion problems. In particular they comprehend in a unifying scheme, free-boundary problems of different nature. We mention specifically problems of fast chemical reaction [5, 8, 9], diffusion in porous media [1, 3, 4, 13, 20, 27], diffusion in porous media of partially saturated gas [14, 25], problems of diffusion involving change of phase of Stefan type [1, 6, 13, 16, 18, 20, 28].

Here we deal with the case in which $\beta(\cdot)$ has a jump at the origin. More precisely we assume $\beta(\cdot)$ is given by

(1.2)
$$\beta(r) = \begin{cases} \beta_1(r) & r > 0 \\ [-v, 0] & r = 0 \\ \beta_2(r) - v & r < 0 \end{cases}$$

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where $\nu > 0$ is a given constant and $\beta_i(\cdot)$ i = 1, 2, are monotone increasing functions in their respective domain of definition, a.e. differentiable and

(1.3)
$$0 < \alpha_0 \le \beta_i^*(r) \le \alpha_1, \quad i = 1, 2$$

for two positive constants α_0 , α_1 .

We introduce some notation and make precise the meaning of solution of (1.1).

Let Ω be a bounded domain in \mathbb{R}^N of boundary $\partial\Omega$ and for $0<\mathbf{T}<\infty$ let $\Omega_{\mathbf{T}}\equiv\Omega\times(0,\mathbf{T}]$, $\Omega(\mathbf{t})\equiv\Omega\times\{\mathbf{t}\}$, $\mathbf{S}_{\mathbf{T}}=\bigcup_{0<\mathbf{t}<\mathbf{T}}\partial\Omega\times\{\mathbf{t}\}$, $\Gamma=\mathbf{S}_{\mathbf{T}}\cup\Omega(0)$.

For q, r \geq l we denote by $L_{q,r}(\Omega_T)$ the Banach space of those measurable functions mapping $\Omega_T \to \mathbb{R}$, with norm defined by

$$\|\mathbf{u}\|_{\mathbf{q},\mathbf{r},\Omega_{m}}^{\mathbf{r}} = \int_{0}^{\mathbf{T}} \|\mathbf{u}\|_{\mathbf{q},\Omega}^{\mathbf{r}}$$
 (t) dt

where

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$$\|\mathbf{u}\|_{\mathbf{q},\Omega}^{\mathbf{q}}$$
 (t) = $\int_{\Omega} |\mathbf{u}(\mathbf{x},t)|^{\mathbf{q}} d\mathbf{x}$.

When q = r = 2, $L_{2,2}(\Omega_T)$ coincides with the Hilbert space $L_2(\Omega_T)$ whose inner product $(\cdot, \cdot)_{2,\Omega_T}$ generates the norm $\|\cdot\|_{2,\Omega_T} \equiv \|\cdot\|_{2,2,\Omega_T}$.

Let $W_2^{1,0}(\Omega_{\underline{T}})$ denote the Hilbert space with inner product

$$(\mathbf{u}, \mathbf{v}) = (\mathbf{u}, \mathbf{v})_{2,\Omega_{\mathbf{T}}} + \sum_{i=1}^{N} (\frac{\partial \mathbf{u}}{\partial \mathbf{x}_{i}}, \frac{\partial \mathbf{v}}{\partial \mathbf{x}_{i}})_{2,\Omega_{\mathbf{T}}}$$

while $\mathbf{W}_{2}^{1,1}(\mathbf{Q}_{T})$ denotes the Hilbert space with inner product

$$(\mathbf{u}, \mathbf{v})_{\mathbf{W}_{2}^{1,1}(\Omega_{\mathbf{T}})}^{1} = (\mathbf{u}, \mathbf{v})_{\mathbf{W}_{2}^{1,0}(\Omega_{\mathbf{T}})}^{1} + (\frac{\partial \mathbf{u}}{\partial \mathbf{t}}, \frac{\partial \mathbf{v}}{\partial \mathbf{t}})_{2,\Omega_{\mathbf{T}}}^{1}$$

Here $\frac{\partial u}{\partial x_1}$, $\frac{\partial u}{\partial t}$ denote generalized derivatives. With $w_2^{1,1}(\Omega_T)$ we denote the space of those elements in $w_2^{1,1}(\Omega_T)$ whose trace on $\partial \Omega \times (0,T]$ is zero.

Let $V_2^{1,0}(\Omega_T)$ denote the Banach space of functions such that the map $t \to u(\cdot,t)$ is continuous with respect to $\|\cdot\|_{2,\Omega}$, and the norm is given by

$$|\mathbf{u}|_{\mathbf{V}_{2}^{1,0}(\Omega_{\mathbf{T}})}^{2} = \sup_{0 \le \mathbf{t} \le \mathbf{T}} \|\mathbf{u}(\cdot,\mathbf{t})\|_{2,\Omega}^{2} + \|\nabla_{\mathbf{x}}\mathbf{u}\|_{2,\Omega_{\mathbf{T}}}^{2} ,$$

where

$$\|\nabla_{\mathbf{x}}\mathbf{u}\|_{2,\Omega_{\mathbf{T}}}^{2} = \sum_{i=1}^{N} \left(\frac{\partial \mathbf{u}}{\partial \mathbf{x}_{i}}, \frac{\partial \mathbf{u}}{\partial \mathbf{x}_{i}}\right)_{2,\Omega_{\mathbf{T}}}$$

From (1.2) it follows that $r \to \beta(r)$ is a relation in $\mathbb{R} \times \mathbb{R}$, whose inverse $\beta^{-1}(\cdot)$ is a function.

<u>Definition</u>: By a weak solution of (1.1) in Ω_T we mean a function $u \in W_2^{1,1}(\Omega_T)$ defined by

$$u \equiv \beta^{-1}(w)$$
,

where w is a function defined in $\Omega_{\underline{\mathbf{T}}}$ such that

$$w \subset \beta(u)$$
,

the inclusion being intended in the sense of the graphs, and w and u satisfy

(1.4)
$$\int_{\Omega} w(\mathbf{x},\tau) \varphi(\mathbf{x},\tau) d\mathbf{x} \Big|_{t_{0}}^{t} + \int_{t_{0}}^{t} \int_{\Omega} \{-w(\mathbf{x},\tau) \frac{\partial}{\partial t} \varphi(\mathbf{x},\tau) + \frac{\partial}{\partial t} (\mathbf{x},\tau,u,\nabla_{\mathbf{x}} u) \cdot \nabla_{\mathbf{x}} \varphi + \frac{\partial}{\partial t} (\mathbf{x},\tau,u,\nabla_{\mathbf{x}} u) \varphi \} d\mathbf{x} d\tau = 0$$

for all $\varphi \in \mathbb{W}_{2}^{1,1}(\Omega_{T})$, and all intervals $[t_{0},t] \in (0,T]$.

If $u \in V_2^{1,0}(\Omega_T)$ is solution of a boundary value problem associated with (1.1), then it satisfies (1.4), the boundary conditions being specified separately. We remark that if in (1.4) we want to allow $t_0 = 0$, then along with $u(x,0) = u_0(x)$, the selection $w_0(x) \in \beta(u_0(x))$ must be given. A common device consists in prescribing $u_0(x) \neq 0$ a.e. in Ω so that $\beta(u_0(x))$ is unambiguously a.e. defined in Ω .

We are not concerned here with the existence of weak solutions of (1.1), for which we refer to [1, 5, 6, 16, 18, 20]. Our results are local in nature and descent only from identity (1.4), so that we need not associate (1.1) with a particular boundary value problem.

Our goal is to prove that a weak solution of (1.1) is continuous in $\,\Omega_{\rm T}^{}.$ For this we introduce the auxiliary function

$$v(x,t) = \beta_0(u(x,t)) = \begin{cases} \beta_1(u(x,t)), & \text{on } [u > 0] \\ 0, & \text{on } [u = 0] \\ \beta_2(u(x,t)), & \text{on } [u < 0] \end{cases}$$

and set

$$w(x,t) = v(x,t) - v(x,t)\chi[v \le 0] ,$$

where $\forall (x,t) \geq 0$ is given by

$$v(x,t) = \begin{cases} v & , & (x,t) \in [v < 0] \\ -w(x,t) & , & (x,t) \in [v = 0] \end{cases},$$

and $\chi(\Sigma)$ denotes the characteristic function of the set Σ .

By virtue of (1.3), if $u \in W_2^{1,1}(\Omega_T)$ then also $v \in W_2^{1,1}(\Omega_T)$, and it will be enough to show the continuity of v in Ω_T .

Setting

$$\vec{a}(\mathbf{x}, \mathbf{t}, \mathbf{v}, \nabla_{\mathbf{x}} \mathbf{v}) = \vec{a}(\mathbf{x}, \mathbf{t}, \beta_0^{-1}(\mathbf{v}), \nabla_{\mathbf{x}} \beta_0^{-1}(\mathbf{v}))$$

$$\mathbf{b}(\mathbf{x}, \mathbf{t}, \mathbf{v}, \nabla_{\mathbf{y}} \mathbf{v}) = \mathbf{b}(\mathbf{x}, \mathbf{t}, \beta_0^{-1}(\mathbf{v}), \nabla_{\mathbf{y}} \beta_0^{-1}(\mathbf{v})) ,$$

identity (1.4) can be rewritten as

$$(1.5) \int_{\Omega} (\mathbf{v}(\mathbf{x},\tau) - \mathbf{v}(\mathbf{x},\tau) \chi [\mathbf{v} \leq 0]) \varphi(\mathbf{x},\tau) d\mathbf{x} \Big|_{t_{0}}^{t} + \int_{0}^{t} \int_{\Omega} \left\{ -(\mathbf{v}(\mathbf{x},\tau) - \mathbf{v}(\mathbf{x},\tau) \chi [\mathbf{v} \leq 0]) \cdot \frac{\partial \varphi}{\partial t} + \stackrel{\rightarrow}{\mathbf{a}}(\mathbf{x},\tau,\mathbf{v},\nabla_{\mathbf{x}}\mathbf{v}) \cdot \nabla_{\mathbf{x}}\varphi + \mathbf{b}(\mathbf{x},\tau,\mathbf{v},\nabla_{\mathbf{x}}\mathbf{v})\varphi \right\} d\mathbf{x} d\tau = 0$$

 $\Psi \in \mathbb{V}_{2}^{1,1}(\Omega_{T})$ and all intervals $[t_{0},t] \in (0,T]$.

The above can be viewed as the weak formulation of

(1.6)
$$\frac{\partial}{\partial t} \beta(v) - \operatorname{div} \overrightarrow{a}(x,t,v,\nabla_{x}v) + b(x,t,v,\nabla_{x}v) \ni 0 \text{ in } \mathcal{D}'(\Omega_{T})$$

where $\beta(\cdot)$ is the maximal monotone graph

(1.7)
$$\beta(r) = \begin{cases} r & r > 0 \\ [0,-v] & r = 0 \\ r - v & r < 0 \end{cases}$$

In what follows we will assume $\beta(\cdot)$ is given as in (1.7).

Throughout the paper we will make the following assumptions on the coefficients $\vec{a} = (a_1, a_2, \dots, a_N)$ and b.

$$[A_1]$$
 $a_i, b \in C[\overline{\Omega}_T \times \mathbb{R}^{N+1}]$, $i = 1, 2, ..., N$.

$$\begin{bmatrix} A_2 \end{bmatrix} \qquad \qquad \sum_{i=1}^{N} a_i(x,t,v,\vec{p}) p_i \ge C_0(|v|) |\vec{p}|^2 - \varphi_0(x,t)$$

 $\begin{aligned} |a_{i}(x,t,v,\vec{p})| &\leq \mu_{0}(|v|)|\vec{p}| + \varphi_{1}(x,t) , & i = 1, 2, ..., N . \\ |b(x,t,v,\vec{p})| &\leq \mu_{1}(|v|)|\vec{p}|^{2} + \varphi_{2}(x,t) , \end{aligned}$

where $C_0(\cdot): \mathbb{R}^+ \to \mathbb{R}^+$ is continuous, decreasing, and strictly positive $\mu_i(\cdot): \mathbb{R}^+ \to \mathbb{R}^+$ are continuous and increasing, i = 0, 1,

and the φ_i , i = 0, 1, 2 are non-negative and satisfy

$$\|\varphi_0, \varphi_2\|_{\hat{\mathbf{q}}, \hat{\mathbf{r}}, \Omega_{\mathbf{T}}}$$
, $\|\varphi_1\|_{2\hat{\mathbf{q}}, 2\hat{\mathbf{r}}, \Omega_{\mathbf{T}}} \leq \mu_2$.

Here μ_2 is a given constant and \hat{q} , \hat{r} are positive numbers linked by

$$\frac{1}{\hat{r}} + \frac{N}{2\hat{q}} = 1 - \kappa_1$$

$$\hat{q} \in \left[\frac{N}{2(1 - \kappa_1)}, \infty \right], \hat{r} \in \left[\frac{1}{1 - \kappa_1}, \infty \right], 0 < \kappa_1 < 1, \text{ for } N \ge 2$$

$$\hat{q} \in (1, \infty), \hat{r} \in \left[\frac{1}{1 - \kappa_1}, \frac{1}{1 - 2\kappa_1} \right], 0 < \kappa_1 < \frac{1}{2}, \text{ for } N = 1$$

We can now state our main result.

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Theorem 1: Let $[A_1]$ - $[A_2]$ hold. Then every essentially bounded weak solution u of (1.1) is continuous in Ω_T .

If (1.1) is associated with an initial boundary value problem of Dirichlet or Neumann type then under suitable assumptions on the boundary conditions the continuity of u can be extended to the closure of $\Omega_{\mathbf{T}}$. For the precise statement of these results we refer to Section 5 Theorem 5.1, 5.2 and 5.3.

Remarks: (i) By the local nature of our arguments, in Theorem 1, the function unneed not be defined in a cylindrical domain, since we can always reduce to this

case by selecting in (1.4) test functions supported in cylindrical domains. Hence for the purpose of proving Theorem 1 one need only to assume that u is, locally essentially bounded in Q and that $u \in \mathbb{W}^{1,1}_{2,loc}(Q)$.

- (ii) It is of interest to know if Theorem 1 holds under the assumptions that u is essentially bounded in Ω_T and $u \in V_2^{1,0}(\Omega_T)$. A step in this direction can be found in Section 6.
- (iii) Assumptions $[A_1]$ $[A_2]$ are the same to those imposed in [18] to study the Hölder continuity of weak solutions of (1.1) with $\beta(r) = r$. In this connection in [22, 23] it is observed that the order of summability \hat{p} , \hat{r} are optimal.

If $r \rightarrow \beta(r)$ is a monotone a.e. differentiable function satisfying (1.3), then the local Hölder continuity of the solution follows from the results of [18]. See also [8, 9] for the corresponding free-boundary problems.

We briefly comment on the regularity results at our knowledge available when $\beta(\cdot)$ is monotone and singular or degenerate.

For N = 1 and b \equiv 0, Fasano, Primicerio and Kamin showed in {15} that, under suitable assumptions on $\vec{a}(x,t,u,\nabla_x u)$, a generalized solution of (1.1) is locally Lipschitz-continuous in Ω_T . Hölder estimates where obtained by Cannon, Henry, Kotlov [10].

In [14] a similar result is obtained for a degenerate $\beta(.)$ of the form

$$\beta(r) = \begin{cases} \beta_1(r) & r < 0 \\ 0 & r \ge 0 \end{cases}$$

where $\beta_1(\cdot)$ satisfies (1.3).

For $N \ge 1$, Caffarelli and Friedman [3] proved the continuity of nonnegative weak solutions of

$$\frac{\partial}{\partial t} u^{\alpha} - \Delta u = 0$$
 , $0 < \alpha \le 1$.

This result has been improved to the Hölder continuity by the same authors in [4].

Recently Caffarelli and Evans [2] have shown that weak solutions of

$$\frac{\partial}{\partial t} \beta(u) - \Delta u \neq 0$$
 in Ω_T

for $\beta(\cdot)$ given by

$$\beta(r) = \begin{cases} \beta_1 r & r > 0 \\ [-v,0] & r = 0 \\ \beta_2 r & r < 0 \end{cases}$$

 β_i , i = 1, 2 positive constants, are continuous. Their method of proof relies strongly on the properties of the Laplacian operator and the absence of lower order terms.

Our approach is completely different from the one in [2], and it is a natural continuation of ideas exposed in [12]. The method consists of a suitable modification of the parabolic version of the De Giorgi's estimates, as appearing in Ladijzenskaja-Solonnikov-Ural'tzeva [18].

The main idea of the proof can be described somehow euristically as follows. The function $(\mathbf{x},\mathbf{t}) \to \mathbf{u}(\mathbf{x},\mathbf{t})$ can be modified in a set of measure zero to yield a continuous representative out of the equivalence class $\mathbf{u} \in \mathbb{W}_2^{1,1}(\Omega_{\mathbf{T}})$ if for every $(\mathbf{x}_0,\mathbf{t}_0) \in \Omega_{\mathbf{T}}$ there exists a family of nested and shrinking cylinders $Q_n(\mathbf{x}_0,\mathbf{t}_0)$ around $(\mathbf{x}_0,\mathbf{t}_0)$, such that the essential oscillation ω_n of \mathbf{u} in $Q_n(\mathbf{x}_0,\mathbf{t}_0)$, tends to zero as $n\to\infty$ in a way determined by the operator in (1.1) and the data.

The statement that a certain quantity, or function, depends upon the data will mean that it can be determined in terms of N, $C_0(\cdot)$, $\kappa_0(\cdot)$, $\kappa_1(\cdot)$, ε_i , $i=0,1,2,-\hat{g}$, \hat{r} , κ_1 , the jump ν of $\beta(\cdot)$ and the essential bound of u over Ω_m .

The paper is organized as follows. Section 2 contains some preliminary material and the derivation of a system of integral inequalities which will be the main tool in the proof of the theorem. Sections 3 and 4 are devoted to the proof of Theorem 1. The continuity up to the boundary is discussed in Section 5.

Finally in section 6 we show that if $u\in V_2^{1,0}(\Omega_T)$ is a weak solution of (1.1) which can be obtained as weak $V_2^{1,0}(\Omega_T)$ -limit of certain approximations of (1.1) (in a sense to be made precise) then in fact the convergence takes place in the topology of the uniform convergence over compacts of Ω_T .

Since the arguments are technically heavy and the symbolism is quite complicated, an effort has been made to render the paper as self-contained as possible.

In view of this we have reproduced certain calculations already known from the literature.

I would like to thank M. Crandall for several helpful discussions on the subject.

2. Preliminary material and integral inequalities:

This section is devoted to the derivation of a system of integral inequalities which will be the main tool in the proof of Theorem 1.

Let $v \in L_{q,r}(\Omega_T)$ and $k \in \mathbb{R}$. Set

$$(v - k)^+ = \max\{(v - k); 0\}; (v - k)^- = \max\{-(v - k); 0\}$$
.

It is obvious that $(v-k)^{\frac{1}{-}} \in L_{q,r}(\Omega_T)$ and it is known that if $v \in W_2^{1,1}(\Omega_T)$ so does $(v-k)^{\frac{1}{-}}$, (see [19]).

With B(R) we denote a ball of radius R in ${\rm I\!R}^N$ and if ${\bf x} \to {\bf v}({\bf x})$ is defined in Ω , and B(R) $\subset \Omega$ we set

$$A_{k,R}^{+} = \{x \in B(R) | v(x) > k\}$$

$$A_{k,R} = \{x \in B(R) | v(x) < k\}$$
.

Also let κ_N denote the measure of the surface of the unit sphere so that meas B(R) = κ_N^{N} .

From now on $(x,t) \rightarrow v(x,t)$ will denote a weak solution of (1.6), and M is a positive real number such that

ess sup
$$|\mathbf{v}| \leq M$$
 .

We will think of $(x,t) \to v(x,t)$ as an arbitrarily selected and fixed representative out of the equivalence class v, so that the map $(x,t) \to v(x,t) \in \mathbb{R}$ is well defined $\forall (x,t) \in \Omega_T$.

We will derive a system of inequalities for v by making particular selections of the test function φ in the identity (1.5).

First we observe that since $v \in W_2^{1,1}(\Omega_T)$, (1.5) can be rewritten as

$$(2.1) - \int_{\Omega} v(\mathbf{x}, \tau) \chi[\mathbf{v} \leq 0] \varphi \, d\mathbf{x} \Big|_{t_{0}}^{t} + \int_{t_{0}}^{t} \int_{\Omega} v(\mathbf{x}, \tau) \chi[\mathbf{v} \leq 0] \, \frac{\partial}{\partial t} \varphi \, d\mathbf{x} d\tau + \int_{t_{0}}^{t} \int_{\Omega} \left\{ \frac{\partial}{\partial t} \, \mathbf{v} \, \varphi + \frac{1}{a}(\mathbf{x}, \tau, \mathbf{v}, \nabla_{\mathbf{x}} \mathbf{v}) \cdot \nabla_{\mathbf{x}} \varphi + b(\mathbf{x}, \tau, \mathbf{v}, \nabla_{\mathbf{x}} \mathbf{v}) \varphi \right\} d\mathbf{x} d\tau = 0$$

 $\Psi_{\varepsilon} \in \mathbb{W}_{2}^{1,1}(\Omega_{\mathbf{T}})$, and any interval $[t_{0},t] \subset (0,T]$.

Next we construct the test functions in (2.1).

Let σ_1 , $\sigma_2 \in (0,1)$ and consider the concentric balls B(R) and B(R - σ_1 R), and the cylinders $Q(R,\lambda) \equiv B(R) \times [t_0,t_0+\lambda]$ and $Q(R-\sigma_1R,\lambda-\sigma_2\lambda) \equiv B(R-\sigma_1R) \times [t_0+\sigma_2\lambda,t_0+\lambda]$, $\lambda>0$.

Define cutoff functions in $Q(R,\lambda)$ as follows

- (a) $\zeta \in C^{\infty}[Q(R,\lambda)]$ such that $\zeta(x,t)|_{\partial B(R)} = 0$ $\forall t \in [t_0,t_0+\lambda]$, $\zeta(x,t_0) = 0$ $\forall x \in B(R)$ and $\zeta(x,t) = 1$, $(x,t) \in Q(R \sigma_1 R, \lambda \sigma_2 \lambda)$, $\frac{\partial}{\partial t} \zeta \geq 0$, $|\nabla_x \zeta| \leq (\sigma_1 R)^{-1}$; $|\frac{\partial}{\partial t} \zeta| \leq (\sigma_2 \lambda)^{-1}$.
 - (b) $\zeta \in C_0^{\infty}(B(R))$ such that $\zeta(x) = 1$, $x \in B(R \sigma_1 R)$, $|\nabla \zeta| \leq (\sigma_1 R)^{-1}$.

For any cylinder $Q(R,\lambda)\subset\Omega_{\widetilde{T}}$ we make the following selections of test function in (2.1)

$$\varphi = \pm (\mathbf{v} - \mathbf{k}) + \zeta^2$$

where $k \in \mathbb{R}$ satisfies

(2.2) ess sup
$$(v - k)^{+} \leq \delta$$
 $Q(R, \lambda)$

for some $\delta > 0$ to be selected, and $(x,t) \to \zeta(x,t)$ is either as in (a) or as in (b).

For simplicity of notation we set

$$-\int_{\Omega} v(\mathbf{x},\tau) \chi[v \le 0] \left[\pm (\mathbf{v} - \mathbf{k}) \pm 1 \zeta^{2} d\mathbf{x} \right]_{t_{0}}^{t} + \int_{t_{0}}^{t} \int_{\Omega} v(\mathbf{x},\tau) \chi[v \le 0] \frac{\partial}{\partial t} \left[\pm (\mathbf{v} - \mathbf{k}) \pm \zeta^{2} \right] d\mathbf{x} d\tau$$

$$= -\phi^{\pm}(\mathbf{k},t_{0},t,\zeta) \qquad , \quad t \in (t_{0},t_{0}+\lambda)$$

and transform and estimate the remaining parts of (2.1) as follows

$$I = \int_{t_0}^{t} \int_{\Omega} \frac{1}{t} \frac{1}{t} v(v - k) + \zeta^2 dx d\tau = \frac{1}{2} \int_{t_0}^{t} \int_{\Omega} \frac{1}{t} \left[(v - k) + \frac{1}{2} z^2 dx d\tau \right] dx d\tau =$$

$$= \frac{1}{2} \left\| (v - k) + \zeta \right\|_{2,\Omega}^{2} (\tau) \left| t - \int_{t_0}^{t} \int_{\Omega} \left[(v - k) + \frac{1}{2} z^2 dx d\tau \right] dx d\tau.$$

To estimate the last two terms in (2.1) we take in account the assumptions $[A_1]$ - $[A_2]$. We have

$$\begin{split} J_{1} &= \int_{t_{0}}^{t} \int_{\Omega} \vec{a}(x,\tau,v,\nabla_{x}v) \cdot \nabla_{x} [\pm (v-k)^{\frac{1}{2}} z^{2}] dx d\tau = \\ &= \sum_{i=1}^{N} \int_{t_{0}}^{t} \int_{\Omega} a_{i}(x,\tau,v,\nabla_{x}v) \frac{\partial}{\partial x_{i}} [\pm (v-k)^{\frac{1}{2}}] z^{2} dx d\tau + \\ &+ \int_{t_{0}}^{t} \int_{\Omega} \vec{a}(x,\tau,v,\nabla_{x}v) [\pm (v-k)^{\frac{1}{2}}] \cdot \nabla_{x} z^{2} dx d\tau \geq \\ &\geq \int_{t_{0}}^{t} \int_{\Omega} c_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}}|^{2} z^{2} dx d\tau - \int_{t_{0}}^{t} \int_{\Omega} c_{0} z^{2} x [(v-k)^{\frac{1}{2}} > 0] dx d\tau \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}}|(v-k)^{\frac{1}{2}} z^{2} dx d\tau - \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} \int_{\Omega} a_{0}(|v|) |\nabla_{x}(v-k)^{\frac{1}{2}} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t_{0}}^{t} |\nabla_{x} z| dx d\tau - \\ &- 2 \int_{t$$

$$J_{2} = -\int_{t_{0}}^{t} \int_{\mathbb{R}^{N}} b(x, \tau, v, v_{x}v) \left[\frac{1}{2} (v - k)^{\frac{1}{2}} \right] \zeta^{2} dx d\tau \ge -\int_{t_{0}}^{t} \int_{\mathbb{R}^{N}} |u_{1}(|v|) |T(v - k)^{\frac{1}{2}}|^{2} .$$

$$(v - k)^{\frac{1}{2}} \zeta^{2} dxd\tau - \int_{t_{0}}^{t} \int_{\Omega} \varphi_{2} (v - k)^{\frac{1}{2}} \zeta^{2} dxd\tau$$

Since ess sup $|v| \le M$, from the assumptions on $C_0(\cdot)$ and $c_1(\cdot)$ we see that T

 $C_0(|v|) \ge C_0(M)$, $\mu_i(|v|) \le \mu_i(M)$, i = 0, 1. From the Cauchy inequality $2ab \le \epsilon a^2 + \epsilon^{-1}b^2$ we have

$$2\int_{t_0}^{t}\int_{\Omega}\mu_0(M)|\nabla_{\mathbf{x}}(\mathbf{v}-\mathbf{k})^{\frac{1}{2}}|(\mathbf{v}-\mathbf{k})^{\frac{1}{2}}|\zeta|^{\frac{1}{2}}\nabla_{\mathbf{x}}|d\mathbf{x}dt| \leq$$

$$\leq \varepsilon \int_{t_0}^{t} \int_{\Omega} \left| \nabla_{\mathbf{x}} (\mathbf{v} - \mathbf{k})^{\frac{1}{2}} \right|^2 |\xi|^2 |\mathrm{d}\mathbf{x} \mathrm{d}\tau| + \varepsilon^{-1} \mu_0^2 (\mathbf{M}) \int_{t_0}^{t} \int_{\Omega} \left[(\mathbf{v} - \mathbf{k})^{\frac{1}{2}} \right]^2 \left| \nabla_{\mathbf{x}} \zeta \right|^2 \mathrm{d}\mathbf{x} \mathrm{d}z$$

and

$$2\int_{t_0}^{t} \int_{\Omega} \varphi_1 \left(v - k\right)^{\frac{1}{2}} \left|\nabla_{\mathbf{x}} \xi\right| d\mathbf{x} d\tau \leq \int_{t_0}^{t} \int_{\Omega} \left[\left(v - k\right)^{\frac{1}{2}}\right]^2 \left|\nabla_{\mathbf{x}} \xi\right|^2 d\mathbf{x} dt + \int_{t_0}^{t} \int_{\Omega} \varphi_1^2 |\xi|^2 |\xi|^2 |\xi|^2 |\xi|^2 |\xi|^2 d\mathbf{x} dt$$

In estimating the integrals in J_2 we recall (2.2), so that combining the estimates for J_1 and J_2 we obtain

$$\begin{split} J_1 + J_2 & \geq (c_0(M) - \epsilon - \delta h_1(M)) \int_{t_0}^{t} \int_{\mathbb{R}} |v_x(v - k)^{\frac{1}{2}}|^2 h^2 dx dt - \\ & - (\epsilon^{-1} h_0^2(M) + 1) \int_{t_0}^{t} \int_{\mathbb{R}} |(v - k)^{\frac{1}{2}}|^2 |v_x(k)|^2 dx dt - \\ & - \int_{t_0}^{t} \int_{\mathbb{R}} |(v_0 + \delta v_2 + v_1^2) \varsigma^2 \chi[(v - k)^{\frac{1}{2}} \neq 0] dx dt \end{split}$$

Next we estimate the last integral above in terms of the measure of the set $[(v-k)^{\frac{1}{2}}>0], \ \ \text{by employing the assumption } [A_2]. \ \ \text{We set}$

$$A_{k,R}^{+}(\tau) = \{x \in B(R) \mid (v - k)^{+}(x,\tau) > 0\}$$
.

Then by the Hölder inequality

$$J^{*} = \int_{t_{0}}^{t} \int_{\Omega} [\varphi_{0} + \delta\varphi_{2} + \varphi_{1}^{2}]_{\zeta}^{2} \chi[(v - k)^{+} > 0] dx d_{T} \le$$

$$\leq \max[1,\delta] \|\varphi_0 + \varphi_2 + \varphi_1^2\|_{\hat{\mathbf{q}},\hat{\mathbf{r}},\cap_{\mathbf{T}}} \left\{ \int_{t_0}^{t} \left[\max_{\mathbf{k},\mathbf{k}} \mathbf{A}_{\mathbf{k},\mathbf{k}}^{+}(\tau) \right] \frac{\hat{\mathbf{q}}-1}{\hat{\mathbf{q}}} \frac{\hat{\mathbf{r}}}{\hat{\mathbf{r}}-1} d\tau \right\} \frac{\hat{\mathbf{r}}-1}{\hat{\mathbf{r}}}.$$

Setting

(2.3)
$$q = \frac{2\hat{q}(1+\kappa)}{\hat{q}-1}$$
, $r = \frac{2\hat{r}(1+\kappa)}{\hat{r}-1}$, $\kappa = \frac{2\kappa}{N}$,

$$J^{\star} \leq \max[1,\delta] \|\varphi_0 + \varphi_2 + \varphi_1^2\|_{\hat{q},\hat{r},\Omega_T} \left\{ \int_{t_0}^{t} \left[\max_{q \in A_{k,R}^+(\tau)} A_{k,R}^+(\tau) \right]^{\frac{r}{q}} d\tau \right\}^{\frac{2}{r}(1+r)}.$$

Since
$$\frac{1}{\hat{r}} + \frac{N}{2q} = 1 + \kappa_1$$
 we have

$$\frac{1}{r} + \frac{N}{2q} = \frac{N}{4} ,$$

and the admissible range of r and q is

$$\begin{cases}
q \cdot (2, \frac{2N}{N-2}], & r \in [2, \infty) \text{ for } N \ge 3 \\
q \cdot (2, \infty), & r \cdot (2, \infty) \text{ for } N = 2 \\
q \cdot (2, \infty), & r \cdot [4, \infty), & \text{for } N = 1
\end{cases}$$

In the estimate of $J_1 + J_2$ we choose

(2.6)
$$\varepsilon = \frac{C_0(M)}{4} \qquad \delta = \min \left\{ \frac{C_0(M)}{4\nu_1(M)}, 1 \right\}$$

so that collecting all the previous estimates we obtain the inequalities

$$\| (v - k)^{\frac{1}{2}} \zeta \|_{2,\Omega}^{2} (t) + \int_{t_{0}}^{t} \int_{\Omega} |\nabla_{\mathbf{x}} (v - k)^{\frac{1}{2}}|^{2} \zeta^{2} dx d\tau \le$$

$$\leq \| (\mathbf{v} - \mathbf{k})^{+} \zeta \|_{2,\Omega}^{2} (\mathbf{t}_{0}) + \gamma \int_{\mathbf{t}_{0}}^{\mathbf{t}_{0}+1} \int_{\Omega} [(\mathbf{v} - \mathbf{k})^{+}]^{2} (|\nabla_{\mathbf{x}} \zeta|^{2} + \zeta |\frac{\partial}{\partial \mathbf{t}} \zeta|) d\mathbf{x} d\tau +$$

$$+ \gamma \left\{ \int_{t_0}^{t_0+\lambda} \left[\text{meas } A_{k,R}^{\frac{1}{2}}(\tau) \right]^{\frac{r}{q}} d\tau \right\}^{\frac{2}{r}} (1+\kappa) + \Phi^{\frac{1}{r}}(k,t_0,t,\zeta), \forall t \in [t_0,t_0^{+\lambda}],$$

where γ is a constant depending only upon the data.

These inequalities are valid for every cylinder $Q(R,\lambda)\subset\Omega_{\overline{L}}$ and every $k\in\mathbb{R}$, satisfying (2.2) with the choice (2.6) of the parameter δ .

If we select the cutoff function $(x,t) + \zeta(x,t)$ as in (a) we see that there exists a constant γ , dependent only upon the data, such that

$$| (v - k)^{\frac{1}{2}} |^{2} v_{2}^{1,0} [B(R - \sigma_{1}R) \times (t_{0} + \sigma_{2}^{\lambda}, t_{0} + \lambda)] | \leq$$

$$= \gamma [(\sigma_{1}R)^{-2} + (\sigma_{2}^{\lambda})^{-1}] | (v - k)^{\frac{1}{2}} |^{2} + (\sigma_{2}^{\lambda}, t_{0} + \lambda) | + \gamma \left\{ \int_{t_{0}}^{t_{0} + \lambda} [meas A_{k,R}^{\frac{1}{2}}(\tau)]^{\frac{r}{q}} d\tau \right\} |^{\frac{2}{r}} (1+\kappa) +$$

$$+ \sup_{t \in [t_{0}, t_{0} + \lambda]} |_{a}^{\frac{1}{r}} (k, t_{0}, t) |,$$

where $2\frac{1}{a}(k,t_0,t)$ coincides with $2\frac{1}{a}(k,t_0,t,\zeta)$ if $\zeta(x,t)$ is selected as in (a).

Choosing now $x \rightarrow \zeta(x)$ as in (b) we have

$$(2.8) \qquad \| (\mathbf{v} - \mathbf{k})^{\frac{1}{2}} \|_{2, \mathbf{A}_{\mathbf{k}, \mathbf{R} + \sigma_{\mathbf{1}} \mathbf{R}}^{\mathbf{t}}}^{2} + \| \nabla_{\mathbf{x}} (\mathbf{v} - \mathbf{k})^{\frac{1}{2}} \|_{2, \mathcal{Q}(\mathbf{R} - \sigma_{\mathbf{1}} \mathbf{R}, \mathbf{k})}^{2} \leq \\ \leq \| (\mathbf{v} - \mathbf{k})^{\frac{1}{2}} \|_{2, \mathbf{A}_{\mathbf{k}, \mathbf{R}}^{+}}^{2} + \gamma (\sigma_{\mathbf{1}} \mathbf{R})^{-2} \| (\mathbf{v} - \mathbf{k})^{\frac{1}{2}} \|_{2, \mathcal{Q}(\mathbf{R}, \lambda)}^{2} + \\ + \gamma \left\{ \int_{\mathbf{t}_{0}}^{\mathbf{t}_{0} + \lambda} [\text{meas } \mathbf{A}_{\mathbf{k}, \mathbf{R}}^{+} (\mathbf{t})]^{\frac{\mathbf{r}_{\mathbf{q}}}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} (1 + \kappa) \\ + \gamma \left\{ \int_{\mathbf{t}_{0}}^{\mathbf{t}_{0} + \lambda} [\text{meas } \mathbf{A}_{\mathbf{k}, \mathbf{R}}^{+} (\mathbf{r})]^{\frac{\mathbf{r}_{\mathbf{q}}}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} (1 + \kappa) \\ + \gamma \left\{ \int_{\mathbf{t}_{0}}^{\mathbf{t}_{0} + \lambda} [\text{meas } \mathbf{A}_{\mathbf{k}, \mathbf{R}}^{+} (\mathbf{r})]^{\frac{\mathbf{r}_{\mathbf{q}}}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} (1 + \kappa) \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{t}_{0} + \lambda} [\text{meas } \mathbf{A}_{\mathbf{k}, \mathbf{R}}^{+} (\mathbf{r})]^{\frac{\mathbf{r}_{\mathbf{q}}}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} (1 + \kappa) \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{t}_{0} + \lambda} [\text{meas } \mathbf{A}_{\mathbf{k}, \mathbf{R}}^{+} (\mathbf{r})]^{\frac{\mathbf{r}_{\mathbf{q}}}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} (1 + \kappa) \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{t}_{0} + \lambda} [\text{meas } \mathbf{A}_{\mathbf{k}, \mathbf{r}}^{+} (\mathbf{r})]^{\frac{2}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} (1 + \kappa) \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} + \lambda} [\text{meas } \mathbf{A}_{\mathbf{k}, \mathbf{r}}^{+} (\mathbf{r})]^{\frac{2}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} (1 + \kappa) \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} + \lambda} [\text{meas } \mathbf{A}_{\mathbf{k}, \mathbf{r}}^{+} (\mathbf{r})]^{\frac{2}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} (1 + \kappa) \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} + \lambda} [\text{meas } \mathbf{r}_{0} (\mathbf{r})]^{\frac{2}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} d\tau \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} + \lambda} [\text{meas } \mathbf{r}_{0} (\mathbf{r})]^{\frac{2}{\mathbf{q}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} d\tau \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} + \lambda} [\text{meas } \mathbf{r}_{0} (\mathbf{r})]^{\frac{2}{\mathbf{r}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} d\tau \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} + \lambda} [\text{meas } \mathbf{r}_{0} (\mathbf{r})]^{\frac{2}{\mathbf{r}}} d\tau \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} (\mathbf{r})} [\text{meas } \mathbf{r}_{0} (\mathbf{r})]^{\frac{2}{\mathbf{r}}} d\tau \right\}^{\frac{2}{\mathbf{r}}} d\tau \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} (\mathbf{r})} [\text{meas } \mathbf{r}_{0} (\mathbf{r})]^{\frac{2}{\mathbf{r}}} d\tau \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} (\mathbf{r})} [\text{meas } \mathbf{r}_{0} (\mathbf{r})]^{\frac{2}{\mathbf{r}}} d\tau \\ + \gamma \left\{ \int_{\mathbf{r}_{0}}^{\mathbf{r}_{0} (\mathbf{r})} [\text{meas } \mathbf{r}_{0} (\mathbf{r})]^{\frac{2}{\mathbf{r}}} d\tau \\ + \gamma \left\{ \int_{\mathbf{r}_{0}$$

with the obvious definition of $\Phi_b^+(k,t_0,t)$.

Roughly speaking, inequalities (2.7) - (2.8) supply some local control on that part of the graph of v which lies above (below) the hyperplane v = k.

Consider a region $\theta \in \Omega_{\mathbf{T}}$ such that

meas[
$$(x,t) \in \mathcal{O}[v(x,t) \leq 0] = 0$$
.

Then for every cylinder $Q(\rho,\lambda)\subset \mathcal{O}, \ \Phi^{+}(\mathbf{k},\mathbf{t}_{0},\mathbf{t})\equiv 0$, hence by choosing the cutoff function $\zeta(\mathbf{x},\mathbf{t})$ as in (a) and as in (b), we see that the functions $(\mathbf{x},\mathbf{t}) \to (\mathbf{v}-\mathbf{k})^{+}(\mathbf{x},\mathbf{t})$ satisfy inequalities (7.1) - (7.2) of [18] page 110. By virtue of the embedding theorem 7.1 of [18] page 120, this implies that $(\mathbf{x},\mathbf{t}) \to \mathbf{v}(\mathbf{x},\mathbf{t})$ is Hölder continuous in every region $\mathcal{O}^{+}\subset \mathcal{O}$. An analogous argument holds for regions \mathcal{O} such that $\max_{\mathbf{x}\in \mathcal{O}}[\mathbf{x},\mathbf{t})\in \mathcal{O}[\mathbf{v}(\mathbf{x},\mathbf{t})\geq 0]=0$.

Because of the presence of the term $\Phi^+(k,t_0,t)$, we do not expect that inequalities (2.7) - (2.8) imply the continuity of the solution, without additional informations contained in the identity (2.1). This is the role of the next two lemmas.

Let $\theta \in \mathbb{R}^+$ and consider the cylinders $Q(R, \theta R^2) \equiv B(R) \times [t_0, t_0 + \theta R^2]$ and $Q(R - \sigma_1 R, \theta R^2) \equiv B(R - \sigma_1 R) \times [t_0, t_0 + \theta R^2]$.

Lemma 2.1. Let $\zeta(x)$ be a cutoff function in $Q(R, \theta R^2)$ chosen as in (b). There exists a constant $C(M, \theta, v)$ such that

$$\iint\limits_{\mathcal{Q}(R,\theta R^2)} \left| \nabla_{\mathbf{x}} \mathbf{v} \right|^2 \, \varsigma^2(\mathbf{x}) d\mathbf{x} dt \, \leq \frac{C(M,\theta,\nu)}{\sigma_1^2} \, \kappa_N^R R^N \quad .$$

<u>Proof:</u> In (2.1) select the test function $\varphi = e^{\lambda v} \zeta^2(x)$, where $\gamma > 0$ will be chosen later. For all $t \in [t_0, t_0 + \theta R^2]$ we have

$$-\int_{\Omega} v(\mathbf{x},\tau) \chi[\mathbf{v} \leq 0] \varphi \ d\mathbf{x} \left[\frac{t}{t_0} - \int_{t_0}^{t} \int_{\Omega} v(\mathbf{x},\tau) \chi[\mathbf{v} \leq 0] \ \frac{\partial}{\partial t} \ \mathbf{v}^{-} e^{-\frac{1}{2} \mathbf{v}^{-}} z^2(\mathbf{x}) d\mathbf{x} dz \right]$$

$$=-\int_{\Omega} v(\mathbf{x},\tau)\chi[\mathbf{v}\leq 0]e^{-\lambda\mathbf{v}^{T}}\zeta^{2}(\mathbf{x})d\mathbf{x}\left[\frac{t}{t_{0}}+v\int_{0}^{t}\int_{0}^{\infty}\frac{\partial}{\partial t}\left(e^{-\lambda\mathbf{v}^{T}}-1\right)+\zeta^{2}(\mathbf{x})d\mathbf{x}d\mathbf{v}\right]$$

This term, and the term $\int_{-\tilde{t}_0}^{\tilde{t}} \int_{\tilde{t}_0}^{\tilde{t}_0} \int_{\tilde{t}_0}^{\tilde{t}_0} v \, \varphi \, dx d\tau / \tau$, can be easily dominated in terms of $\tilde{C}_{-\tilde{t}_0}R^N$ where \tilde{C} depends upon M, $\tilde{\tau}_0$ and $\tilde{\tau}_0$.

On the other hand using the assumptions $[A_1]$ - $[A_2]$, standard calculations yield

$$\begin{split} \int_{t_0}^t \int_{\mathbb{R}^2} & \{ \tilde{a}(\mathbf{x},\tau,\mathbf{v},\nabla_{\mathbf{x}}\mathbf{v}) \nabla_{\mathbf{x}}\mathbf{v} + b(\mathbf{x},\tau,\mathbf{v},\nabla_{\mathbf{x}}\mathbf{v}) \mathbf{v} \} d\mathbf{x} d\tau \geq \\ & \geq \left[(C_0(M) - \varepsilon - \varepsilon_1(M)) \right] \int_{t_0}^t \int_{\mathbb{R}^2} \left| e^{AV} |\nabla_{\mathbf{x}}\mathbf{v}|^2 |\zeta^2(\mathbf{x}) d\mathbf{x} d\tau - \right. \\ & - \left| e^{AV} \int_{t_0}^t \int_{\mathbb{R}^2} \left((\nabla_{\mathbf{v}} + \varphi_2) \zeta^2(\mathbf{x}) d\mathbf{x} d\tau - 2e^{AW} \int_{t_0}^t \int_{\mathbb{R}^2} |\nabla_{\zeta}| \nabla_{\zeta} |d\mathbf{x} d\tau - \left. - \frac{4}{\varepsilon} |\varepsilon_0(M) e^{AW} \int_{t_0}^t \int_{\mathbb{R}^2} |\nabla_{\zeta}|^2 |d\mathbf{x} d\tau \right]. \end{split}$$

Selecting $\epsilon = \mu_1(M)$ and $\lambda = \frac{4}{C_0(M)} \mu_1(M)$ we conclude that there exists a constant \bar{C} depending upon M such that

$$e^{-tM} = \iint_{\Sigma(R_{+} \cap R^{2})} |\nabla_{\mathbf{x}} \mathbf{v}|^{2} |\zeta^{2}(\mathbf{x}) d\mathbf{x} d\tau| \leq \overline{C} \kappa_{N} R^{N} +$$

$$+ \overline{\overline{c}} \left\{ \int_{t_0}^{t_0 + vR^2} \int_{0}^{\infty} \left[(\varphi_0 + \varphi_2 + \varphi_1^2) \zeta^2(\mathbf{x}) d\mathbf{x} d\tau + \int_{t_0}^{t_0 + \theta R^2} \int_{0}^{\infty} \left| \nabla \zeta \right|^2 d\mathbf{x} d\tau \right\}.$$

We recall that $|\nabla \zeta| < (\sigma_1 R)^{-1}$ and treat the integral involving the φ_1 = 0, 1, 2 as previously, to obtain

$$\iint\limits_{\mathbb{R}^2(\mathbb{R}, |\mathbb{F}^2)} ||\nabla_{\mathbf{x}} v||^2 |\zeta^2(\mathbf{x}) d\mathbf{x} d\tau \leq \frac{C(M, -, \vee)}{\sigma_1^2} |\kappa_N^R|^N \quad .$$

This inequality will be employed to prove the following lemma.

Lemma 2.2: Let $k \in \mathbb{R}^+, \mu \ge \text{ess sup } (v - k)^+$ and $0 < \eta < \mu$ $Q(R, \theta R^2)$

Set

The second secon

$$\psi(\mathbf{x},t) = \ln^{+} \left[\frac{\mu}{\mu - (\mathbf{v} - \mathbf{k})^{+} + \eta} \right] = \max \left\{ \ln \left[\frac{\mu}{\mu - (\mathbf{v} - \mathbf{k})^{+} + \eta} \right]; 0 \right\},$$

then there exists a constant $C = C(\theta)$ such that for all $t \in [t_0, t_0 + \theta R^2]$

$$\int_{B(R-\sigma_1R)} \psi^2(\mathbf{x},t) d\mathbf{x} \leq \int_{B(R)} \psi^2(\mathbf{x},t_0) d\mathbf{x} +$$

$$+\frac{C}{\sigma_1^2}$$
 (1 + ln $\frac{\mu}{\eta}$) (1 + $\frac{R^{N\kappa}}{\eta^2}$) κ_N^{R} .

Remark: For simplicity of notation we will use the same symbol ψ for $\psi(\mathbf{x},t)$ and $\tilde{\psi}(\mathbf{v}(\mathbf{x},t))$. In what follows ψ' will mean $\frac{\partial}{\partial \mathbf{v}}\tilde{\psi}$.

Proof: In (2.1) we select $\varphi = (\psi^2)^{'} \zeta^2(\mathbf{x})$, where $\zeta(\mathbf{x})$ is chosen as in (b). It is apparent that $\varphi \in \mathring{W}_2^{1,1}(\Omega_{\mathbf{T}})$, and that $(\psi^2)^{''} = 2(1+\psi)(\psi^{'})^2$. Since $(\psi^2)^{'}$ vanishes at those points $(\mathbf{x},\mathbf{t}) \in \Omega_{\mathbf{T}}$ where $(\mathbf{v} - \mathbf{k})^+ \leq \eta$, and $\eta > 0$, the terms involving $\forall (\mathbf{x},\mathbf{t})\chi[\mathbf{v} \leq 0]$ in (2.1) does not give any contribution. The term involving $\frac{\partial}{\partial \mathbf{t}} \mathbf{v}$ gives

$$\int_{t_0}^{t} \int_{\Omega} \frac{\partial}{\partial t} \mathbf{v}(\psi^2)' \zeta^2(\mathbf{x}) d\mathbf{x} d\tau = \int_{\Omega} \psi^2(\mathbf{x}, \tau) \zeta^2(\mathbf{x}) d\mathbf{x} \Big|_{t_0}^{t}.$$

We estimate the remaining terms as follows

$$\begin{split} &\int_{t_{0}}^{t} \int_{\mathbb{R}} |\tilde{a}(x,\tau,v,\nabla_{x}v)\{2(1+\phi)(\psi^{*})^{2}|\nabla_{x}v|\xi^{2}(x)+(\psi^{2})^{*}|\nabla\xi^{2}(x)\}dxd\tau \geq \\ &= 2C_{0}(\Omega)\int_{t_{0}}^{t} \int_{\Omega} |(1+\psi)|\nabla_{x}\psi|^{2}|\xi^{2}(x)dxd\tau = \\ &-2\int_{t_{0}}^{t} \int_{\Omega} |\varphi_{0}(x,\tau)(1+\psi)(\psi^{*})^{2}|\xi^{2}(x)dxd\tau = \\ &-4c_{0}(\Omega)\int_{t_{0}}^{t} \int_{\Omega} |\psi^{2}||\nabla_{x}\psi||\xi(x)\psi^{\frac{1}{2}}|\nabla\xi|dxd\tau = \\ &-\int_{t_{0}}^{t} \int_{\Omega} |\varphi_{1}(x,\tau)(\psi^{2})|^{*}|\nabla|\xi^{2}(x)|dxd\tau = \\ &-\int_{t_{0}}^{t} \int_{\Omega} |\varphi_{1}(x,\tau)(\psi^{2})|^{*}|\nabla|\xi^{2}(x)|dxd\tau = \\ &-\int_{t_{0}}^{t} \int_{\Omega} |\{(\xi^{*})^{2}\varphi_{0}(x,\tau)(1+\psi)|\nabla_{x}\psi|^{2}|\xi^{2}(x)dxd\tau = \\ &-\int_{t_{0}}^{t} \int_{\Omega} |\{(\xi^{*})^{2}\varphi_{0}(x,\tau)(1+\psi)\xi^{2}(x)+\varphi_{1}(x,\tau)|\nabla\xi^{2}|(\psi^{2})^{*}\}dxd\tau = \\ &-\int_{t_{0}}^{t} |\{(\xi^{*})^{2}\varphi_{0}(x,\tau)(1+\psi)\xi^{2}(x)+\varphi_{1}(x,\tau)(1+\psi)\xi^{2}(x)+\varphi_{1}(x,\tau)(1+\psi)\xi^{2}(x)+\varphi_{1}(x,\tau)(1+\psi)\xi^{2}(x) + \frac{1}{2}|\nabla\xi^{2}(x)+\varphi_{1}(x,\tau)(1+\psi)\xi^{2}(x)+\varphi_{1}(x,\tau)(1+\psi)\xi^{2}(x) + \frac{1}{2}|\nabla\xi^{2}(x)+\varphi_{1}(x)+\varphi_{1}(x)+\varphi_{1}(x)+\varphi_{1}(x)+\varphi_{1}(x)+\varphi_{1}(x)$$

For the lower order terms we have

$$\int_{t_0}^{t} \int_{\mathbb{R}^{d}} |b(\mathbf{x}, \varepsilon, \mathbf{v}, \nabla_{\mathbf{x}} \mathbf{v}) (\varepsilon^2) |^{t} \varepsilon^2(\mathbf{x}) |d\mathbf{x} d\tau \leq 2 h_1(\mathbf{M}) \int_{t_0}^{t} \int_{\mathbb{R}^{d}} ||\nabla_{\mathbf{x}} \mathbf{v}||^2 ||\xi^2(\mathbf{x}) d\mathbf{x} d\tau + \int_{t_0}^{t} \int_{\mathbb{R}^{d}} ||\varphi_2(\mathbf{x}, \tau)| (\varepsilon^2) |^{t} ||\xi^2(\mathbf{x}) d\mathbf{x} d\tau|.$$

Since
$$|\psi\psi'| |\nabla_{\mathbf{x}} \mathbf{v}|^2 = |\psi^2| |\nabla_{\mathbf{x}} \psi| |\psi^2| |\nabla_{\mathbf{x}} \mathbf{v}|$$
 we have

$$2\mu_{1}(M)$$
 $\int_{t_{0}}^{t} \int_{\Omega} |\psi\psi^{\dagger}| \nabla_{\mathbf{x}} v |^{2} \zeta^{2}(\mathbf{x}) d\mathbf{x} d\tau \leq$

$$\leq \varepsilon \int_{t_0}^{t} \int_{\Omega} (1 + \psi) |\nabla_{\mathbf{x}}\psi|^2 \zeta^2(\mathbf{x}) d\mathbf{x} d\tau +$$

$$+ \varepsilon^{-1} 4\mu_1^2(M) \int_{t_0}^{t} \int_{\Omega} \psi |\nabla_x v|^2 \zeta^2(x) dx d\tau .$$

Collecting the previous estimates gives

$$\int_{\Omega} \psi^{2} \zeta^{2}(\mathbf{x}) d\mathbf{x} \left| \begin{matrix} t \\ t_{0} \end{matrix} + \left[2C_{0}(\mathbf{M}) - 2\varepsilon \right] \int_{t_{0}}^{t} \int_{\Omega} (1 + \psi) \left| \nabla_{\mathbf{x}} \psi \right|^{2} \zeta^{2}(\mathbf{x}) d\mathbf{x} d\tau \le 0$$

$$\leq 2 \int_{t_0}^{t} \int_{\Omega} \{(\psi')^2 (1 + \psi) [\varphi_0 + \varphi_1^2] + \psi \psi' \varphi_2 \} \zeta^2(x) dx d\tau +$$

$$+ \left[2 + \epsilon^{-1} 4(4\mu_0^2(M) + \mu_1^2(M))\right] \int_{t_0}^{t} \int_{\Omega} \psi\{\left|\nabla_{\mathbf{x}} \mathbf{v}\right|^2 \zeta^2(\mathbf{x}) + \left|\nabla_{\mathbf{x}} \zeta\right|^2\} d\mathbf{x} d\tau .$$

We select $\varepsilon = C_0(M)$ and observe that since $n < \mu$, $\psi^* < \frac{1}{\eta}$ and $\psi \le \ln \frac{\mu}{\eta}$. Moreover we recall that $|\nabla \zeta| < (\sigma_1 R)^{-1}$ and lemma 2.1 in estimating the last term in the inequality above. This yields the existence of a constant $\tilde{C}(M, \nu, \theta)$ such that

$$\int_{\mathbb{S}} |\psi^{2}|^{2} |\xi^{2}(\mathbf{x}) d\mathbf{x}| \Big|_{\mathbf{t}_{0}}^{\mathbf{t}} \leq \frac{\tilde{c}}{z_{1}^{2}} (1 + \tilde{c}n \frac{\mu}{\eta}) \left\{ \frac{1}{\eta^{2}} \int_{\mathbf{t}_{0}}^{\mathbf{t}_{0} + \tilde{c}R^{2}} \int_{\mathbb{B}(\mathbb{R})} [\varphi_{0} + \varphi_{2} + \varphi_{1}^{2}] d\mathbf{x} d\mathbf{x} + \kappa_{N} \mathbf{R}^{N} \right\} \leq \frac{\tilde{c}}{z_{1}^{2}} (1 + \tilde{c}n \frac{\pi}{\eta}) \left\{ \frac{1}{\eta^{2}} \|\varphi_{0} + \varphi_{2} + \varphi_{1}^{2}\|_{\hat{\mathbf{q}}, \hat{\mathbf{r}}, \Omega_{\mathbf{T}}} \frac{2}{\mathbf{r}} (1 + \kappa) \frac{2}{\kappa_{N}^{2}} (1 + \kappa)^{-1} \cdot \kappa_{N} \mathbf{R}^{N+N\kappa} + \kappa_{N} \mathbf{R}^{N} \right\}.$$

This proves the lemma.

Remarks: (i) If $k \le 0$ and $\frac{1}{2}$ ess sup $(v - k)^{-}$, then an analogous lemma holds for

$$\overline{\zeta}(\mathbf{x},t) = i \mathbf{n}^{+} \begin{bmatrix} \frac{1}{\sqrt{1 - (\mathbf{v} - \mathbf{k})^{T} + \eta_{1}}} \end{bmatrix} ; 0 < \eta < \overline{\mu} .$$

The proof is the same except for minor changes.

(ii) The proof shows that $C(\theta)$ increases with θ . We will use lemma 2.2 with $0 \le \theta \le 1$ and $C(\theta)$ replaced by C = C(1).

We report a lemma due to De Giorgi [11] which will be used as we proceed. Lemma 2.3 (De Giorgi): Let $\mathbf{v} \times \mathbf{W}_1^1(B(R))$ and let k,ℓ be real numbers such that k > k. Then

(2.9)
$$(\cdot - k) \operatorname{meas} A_{k,R} \leq D \frac{R^{N+1}}{\operatorname{meas} (B(R) \setminus A_{k,R})} \int_{A_{k,R} \setminus A_{k,R}} |\nabla v| dx ,$$

where D is a constant depending only upon the dimension N.

Inequality (2.9) holds for domains other than balls. We refer to [18,19] for details noticing for later use that it is valid for convex domains.

Finally if Q is a cylindrical domain in \mathbb{R}^{N+1} , $\overset{\circ 1,0}{v_2}(\mathbb{Q})$ denotes the subspace of $v_2^{1,0}(\mathbb{Q})$ of functions whose trace is zero on the lateral boundary of Q, equipped with the same norm as $v_2^{1,0}(\mathbb{Q})$.

The proof of the following embedding lemma can be found in [18] page 74-77. Lemma 2.4: If $v \in V_2^{1,0}(Q)$ then $v \in L_{q,r}(Q)$ where q,r satisfy (2.4) - (2.5). Moreover there exists a constant β depending only upon the dimension N such that

(2.10)
$$\|v\|_{q,r,Q} \leq \beta |v|_{v_2^{1,0}(Q)}$$

If q = r = 2 then

(2.11)
$$\|v\|_{2,Q} \leq \beta[\text{meas}[|v| \neq 0] \cap Q]^{\frac{1}{N+2}} |v|_{2,Q}^{1,0}(Q)$$

If $v \in V_2^{1,0}(Q)$ then (2.10) is still valid. Moreover if p=r=2 and if $Q \in \Omega \times (0,T)$

(2.12)
$$\|v\| \le C[\text{meas}[|v| \neq 0] \cap Q]^{\frac{1}{N+2}} |v| v_2^{1,0}(Q)$$

where
$$C = 2\hat{E} + (T^{\frac{N}{2}} \text{ meas } \Omega^{-1})^{\frac{1}{N+2}}$$
.

3. The main proposition:

Throughout this section we let (x_0,t_0) and 0 > 0 and for R > 0, Q_R will denote the cylinder

$$Q_R = \{ | x - x_0 | < R \} \times [t_0 - R^2, t_0]$$
.

Let $R_0 \le \frac{1}{2}$ be so small that $Q_{2R_0} \subseteq \Omega_T$, set

$$\mu^{+} = \text{ess sup } v$$
 ; $\mu^{-} = \text{ess inf } v$, Q_{2R_0}

and denote with ω any positive real number such that

$$2M \ge \omega \ge \text{ess osc } \mathbf{v} = \mu^+ - \mu^-$$
. Q_{2R_0}

For $k \in \mathbb{R}$ and $0 < R \le 2R_0$ we set

$$Q_R^+(k) \in \{(x,t) \in Q_R | v(x,t) > k\}$$

$$Q_R^{\dagger}(k) \equiv \{(x,t) \in Q_R^{\dagger}(v(x,t) < k\}$$
.

Finally we let s denote the smallest positive integer such that

$$\frac{2M}{2^s} < \delta \quad , \quad s \ge 2 \quad ,$$

where δ is the number introduced in (2.6).

The goal of this section is to prove the following result

Proposition 3.1: Let ω be any positive number such that

$$2M \geq \omega \geq \text{ess osc } v$$
 .
$$\frac{Q}{2R_0}$$

Then there exist numbers $s_0 \in \mathbb{N}$, A, a > 1, h > 1, $\xi_{\star} < 1$ such that

ess osc
$$v \le \omega(1 - \frac{1}{s_0^{+A/\omega}a})$$

where $R_* = \xi_* (2R_0)^h$, provided that

$$\frac{\omega}{2^{s_0+A/\omega}a} \geq (2R_0)^{\frac{N\kappa}{2}}.$$

The numbers s_0 , A, a, h, ξ_\star depend uniquely upon the data and not upon R_0 nor ω .

Without loss of generality we may assume that

$$|\mu^-| \leq \mu^+ .$$

If the reverse inequality holds the arguments are similar. Also we will assume that

(3.3)
$$\operatorname{ess \ osc \ v} = \mu^{+} - \mu^{-} > \frac{\omega}{2^{s-1}} ,$$

and treat later the case $\mu^+ - \mu^- \le \frac{\omega}{2^{s-1}}$.

Notice that (3.2) - (3.3) imply that

(3.4)
$$\mu^{+} - \frac{\omega}{2^{s}} > \left| \frac{\omega}{2^{s}} + \mu^{-} \right| \geq 0 .$$

Observe moreover that we may assume

(I)
$$H = \text{ess sup } (v - (\mu^{-} + \frac{\omega}{2^{s}}))^{-} > \frac{\omega}{2^{s+1}}$$
.

Indeed if (") is violated

$$-\underset{\mathcal{Q}_{R_0}}{\text{ess inf }} v \leq -\mu^{-} - \frac{\omega}{2^s} + \frac{\omega}{2^{s+1}}$$

and adding less sup v on the left hand side and $\mu^{\frac{1}{4}}$ on the right hand side we $\frac{Q_{R_0}}{}$

obtain

ess osc
$$v \le \omega(1 - \frac{1}{2^{s+1}})$$

and Proposition 3.1 becomes trivial.

Proposition 3.1 will be a consequence of a series of lemmas which we state and prove independently.

Lemma 3.1: There exists a number c_0 depending only upon the data and independent of ω and R_0 such that if

$$\max Q_{R_0}^{-}(n^{-} + \frac{\omega}{2^s}) \le C_0^{-\omega} + \frac{2\kappa_1}{N+2\kappa_1} \kappa_N^{N+2},$$

then either

(i)
$$|| = \underset{\Sigma_{R_0}}{\text{ess sup}} (v - (L^{-} + \frac{\omega}{2^{S}}))^{-} \le \frac{\frac{N\kappa}{2}}{N_0}$$

CT

(ii)
$$\max_{\frac{R_0}{2}} 2 \frac{1}{2} \left(\frac{1}{2} + \frac{4}{2} - \frac{1}{2} H \right) = 0 .$$

Here x_1 is the number appearing in the assumptions $[\mathbf{A}_1]$ - $[\mathbf{A}_2]$.

Proof of lemma 3.1: Consider inequalities (2.7) for the function $(x,t) \rightarrow (v-k)^{-}$, $\mu^{-} \leq k \leq \mu^{-} + \frac{\omega}{2^{s}}$ in the cylinders Q_{R} , $0 < R \leq R_{0}$. Notice that in view of (3.1)

ess sup
$$(v - k)^{-} \le \frac{\omega}{2^{s}} < \delta$$

so that the use of (2.7) for $\mu \le k \le \mu^- + \frac{\omega}{2^s}$ is justified.

We estimate $\Phi_a^-(k,t_0-R^2,t_0)$ in (2.7) by distinguishing the cases of $k\leq 0$ and k>0.

If $k \le 0$ then $\Phi_a(k, t_0 - R^2, t_0) = 0$.

If k > 0 we have

$$\Phi_{a}^{-}(k,t_{0}-R^{2},t_{0}) = -\int_{\Omega} v(x,\tau)\chi[v \leq 0]v^{-}\zeta^{2}(x,\tau)dx\Big|_{t_{0}-R^{2}}^{t_{0}} -$$

$$-k \int_{\Omega} v(x,\tau) \chi[v \le 0] \zeta^{2}(x,\tau) dx \Big|_{t_{0}-R^{2}}^{t_{0}} + \int_{t_{0}-R^{2}}^{t_{0}} \int_{\Omega} v(x,\tau) \chi[v \le 0] (v - k)^{-1}$$

$$\frac{\partial}{\partial t} \zeta^{2}(\mathbf{x},\tau) d\mathbf{x} d\tau + \int_{t_{0}-R^{2}}^{t_{0}} \int_{\Omega} v \frac{\partial}{\partial t} v^{-} \zeta^{2}(\mathbf{x},\tau) d\mathbf{x} d\tau \leq$$

$$\leq 2\nu \int_{t_0-R^2}^{t_0} \int_{\Omega} (v - k)^{-\frac{\partial}{\partial t}} \zeta(x,\tau) dx d\tau - \nu \int_{t_0-R^2}^{t_0} \int_{\Omega} v^{-}(x,\tau) \frac{\partial}{\partial t} \zeta^{2}(x,\tau) dx d\tau .$$

If $(x,t) \rightarrow \zeta(x,t)$ is selected as in (a), $\zeta_t \ge 0$ so that $v^- \zeta_t \ge 0$ and

$$\phi_{a}^{-}(k,t_{0}-R^{2},t_{0}) \leq \frac{2v}{\sigma_{2}R^{2}} \int_{Q_{R}} (v-k)^{-} dx d\tau$$
.

Inequalities (2.7) now read

$$| (v - k)^{-1} |_{V_{2}^{1}, 0}^{2} \{ B(R - \sigma_{1}R) \times (t_{0} - (1 - \sigma_{2})R^{2}, t_{0}) \} | \leq$$

$$\leq \gamma [(\sigma_{1}R^{-2} + (\sigma_{2}R^{2})^{-1}] \| (v - k)^{-1} \|_{2, Q_{R}}^{2} +$$

$$+ \gamma \left\{ \int_{t_{0} - R^{2}}^{t_{0}} [meas A_{k, R}^{-}(\tau)]^{\frac{r}{q}} d\tau \right\}^{\frac{2}{r}(1 + \kappa)} +$$

$$+ 2v(\sigma_{2}R^{2})^{-1} \int_{Q_{R}}^{} (v - k)^{-1} dx d\tau .$$

Inequalities (3.5) hold for all $\mu \le k \le \mu^- + \frac{\omega}{2^5}$, all σ_1 , $\sigma_2 \in$ (0,1) and all cylinders Q_R , $0 \le R \le R_0$.

Set

$$R_n = \frac{R_0}{2} + \frac{R_0}{2^{n+2}}$$

$$\bar{R}_n = \frac{R_0}{2} + \frac{3R_0}{2^{n+4}}$$
,

and consider the cylinders $Q_{\underset{n}{R}}$ and

$$\tilde{Q}_n = \{ |\mathbf{x} - \mathbf{x}_0| < \tilde{\mathbf{R}}_n \} \times \{ \mathbf{t}_0 - \mathbf{R}_n^2, \mathbf{t}_0 \}$$

$$\overline{\overline{Q}}_n \in \{|\mathbf{x} - \mathbf{x}_0| < \overline{R}_n\} \times \{t_0 - \overline{R}_{n+1}^2, t_0\} .$$

Obviously

$$Q_{R_{n+1}} \times \overline{Q}_n \times \overline{Q}_n \times Q_{R_n}$$
.

Construct smooth cutoff functions $x \to \zeta_n(x)$ as follows

(i)
$$\zeta_n(x) = 1 |x - x_0| < R_{n+1}$$

(ii)
$$\zeta_n(x) = 0$$
 $|x - x_0| > \frac{1}{2} [R_n + R_{n+1}] = \overline{R}_n$

(iii)
$$\left|\nabla\zeta_{n}(x)\right| \leq 2^{n+4}/R_{0}$$
.

For simplicity of notation set

$$k_1 = \mu^- + \frac{\omega}{2^s} .$$

Our purpose is to apply (3.5) to the pair of cylinders $\mathcal{Q}_{\underset{n}{R}}$ and $\overline{\tilde{\mathbb{Q}}}_{n}$ for the decreasing levels

$$k_n = (k_1 - \frac{1}{2} H) + \frac{1}{2^n} H$$
 , $n = 1, 2, ...$

which as easily verified satisfy

$$\mu^- \leq k_n \leq k_1$$
.

Set

$$y_n = \int_{Q_{R_n}} \int [(v - k_n)^T]^2 dx d\tau$$
 and

$$z_{n} = \begin{cases} t_{0} & \frac{r}{r} \\ t_{0} - R_{n}^{2} & [\text{meas } A_{n,R_{n}}^{-}(\tau)]^{\frac{q}{q}} d\tau \end{cases}$$

As n → ∞

$$y_n \rightarrow y_\infty = \int_{\frac{R_0}{2}} [(v - (k_1 - \frac{1}{2} H))^{-}]^2 dxd\tau$$

$$z_{n} \rightarrow z_{\infty} = \left\{ \int_{t_{0}^{-}}^{t_{0}^{-}} \left[\max_{k_{1}^{-}} \frac{A_{k_{1}^{-}}^{-}}{\frac{1}{2} H_{r} R_{0}^{2}} (\tau) \right]^{\frac{r}{q}} d\tau \right\}^{\frac{2}{r}}.$$

Therefore the lemma will be proved if we can show that $y_{\infty} = z_{\infty} = 0$.

Claim: The numbers

$$Y_{n} = \frac{Y_{n}}{H^{2}R_{0}^{N+2}}, Z_{n} = \frac{z_{n}}{R_{0}^{N}},$$

satisfy the recursion inequalities

[I]
$$Y_{n+1} \leq \frac{\tilde{c}2^{5n}}{H} \begin{bmatrix} 1 + \frac{2}{N+2} & \frac{2}{N+2} \\ Y_n & + Y_n^{N+2} & Z_n^{1+\kappa} \end{bmatrix}$$

$$z_{n+1} \leq \frac{\tilde{c}2^{5n}}{H} \left[Y_n + z_n^{1+\kappa} \right] ,$$

where

$$\tilde{C} = 2^{12} \beta^2 \gamma_0 .$$

Here β is the constant appearing in the embedding lemma 2.4, and $\gamma_0 = \max\{\nu, (1+\gamma)(2M+\delta)\}.$

We remark that \tilde{C} depends only upon the data and the dimension N.

Proof of the claim: We use here the method of [18] page 106. Set

$$P_n = \int_{t_0^{-R_{n+1}^2}}^{t_0} [meas A_{k_{n+1},R_{n+1}}^{-}(\tau)] d\tau$$
,

and observe that

$$P_n \le (k_n - k_{n+1})^{-2} y_n = \left(\frac{2^{n+1}}{H}\right)^2 y_n$$
.

By virtue of lemma 2.4 applied over the cylinder $\bar{\bar{Q}}_n$ we have

(3.6)
$$y_{n+1} \leq \int_{\overline{Q}_n} \int [(v - k_{n+1})]^{-2} \zeta_n^2(x) dx d\tau \leq$$

$$\leq \beta^2 p_n^{\frac{2}{N+2}} | (v - k_{n+1})^{-\zeta_n} |_{V_2^{1,0}(\overline{\mathbb{Q}}_n)}^2$$

Estimate of $|(v - k_{n+1})^{-} \zeta_{n}|_{v_{2}^{1,0}(\overline{\overline{Q}}_{n})}^{2}$:

$$\begin{split} & | (v - k_{n+1})^{-\zeta_{n}} |_{V_{2}^{1,0}(\overline{Q}_{n})}^{2} \leq | (v - k_{n+1})^{-1} |_{V_{2}^{1,0}(\overline{Q}_{R_{n}})}^{2} + \\ & + 2 \int_{\overline{Q}_{n}} \int [(v - k_{n+1})^{-1}]^{2} | \nabla \zeta_{n}(x)| dx dt = J_{n}^{(1)} + J_{n}^{(2)} . \end{split}$$

For $J_n^{(2)}$ we have

$$J_{n}^{(2)} \leq 4R_{0}^{-1}2^{n+4} \int_{Q_{R_{N}}} \int \{ [(v - k_{n})^{-}]^{2} + (k_{n} - k_{n+1})^{2} \}_{\chi} (v < k_{n+1}) dx d\tau \leq C_{N}^{(2)}$$

$$\leq 8 2^{n+4} y_n$$
.

In order to estimate $J_n^{(1)}$ we use inequalities (3.5) for the pair of cylinders $\bar{\tilde{Q}}_n$ and \tilde{Q}_R . Notice that in this connection

$$(\sigma_1 R_n)^{-2} = R_0^{-2} 2^{2(n+4)} ; (\sigma_2 R_n^2)^{-1} < R_0^{-2} 2^{n+3}$$
,

so that from (3.5) we deduce

$$J_{n}^{(1)} \leq \frac{2\gamma}{R_{0}^{2}} 2^{2(n+4)} \int_{Q_{R_{n}}} \int \left[(v - k_{n+1})^{2} \right]^{2} dx d\tau + \gamma z_{n}^{1+\kappa} +$$

$$+\frac{v2^{n+4}}{R_0^2}\int_{Q_{R_n}} \int (v-k_{n+1})^{-} dx d\tau \le$$

$$\leq \frac{8\gamma}{R_0^2} 2^{2(n+4)} y_n + \gamma z_n^{1+\kappa} + \frac{\nu}{R_0^2} 2^{n+4} H \int_{t_0-R_n^2}^{t_0} [\text{meas } A_{k_{n+1},R_n}^{-}(\tau)] d\tau .$$

Since

$$\int_{t_0-R_0^2}^{t_0} \left[\text{meas } A_{k_{n+1},R_n}^-(\tau) \right] d\tau \leq (k_n - k_{n+1})^{-2} \ y_n = \frac{(2^{n+1})^2}{H^2} \ y_n \ , \ \text{and}$$

since $H \le 2M + 5$, setting $\gamma_0 = \max\{v, (1 + \gamma)(2M + 5)\}$ above yields

$$J_{n}^{(1)} \leq \frac{8(1+2^{4})\gamma_{0}2^{3(n+1)}}{R_{0}^{2}H} \{y_{n} + R_{0}^{2} z_{n}^{1+\kappa}\}.$$

Combining the estimates for $J_n^{(1)}$ and $J_n^{(2)}$ we have

$$(3.7) \qquad \left| (v - k_{n+1})^{-\zeta_{n}} \right|_{V_{2}^{1}, 0(\overline{\mathbb{Q}}_{n})}^{2} \leq \frac{2^{8} \gamma_{0} (2^{n+1})^{3}}{R_{0}^{2} H} \left\{ y_{n} + R_{0}^{2} z_{n}^{1+\kappa} \right\}$$

Estimate of Y_{n+1} : We carry (3.7) in (3.6) and employ the estimate of P_n to obtain

$$\frac{y_{n+1}}{H^2} \leq \frac{2^{8\beta^2} \gamma_0 (2^{n+1})^{3+\frac{4}{N+2}}}{R_0^2 H} \left\{ \left(\frac{y_n}{H^2} \right)^{1+\frac{2}{N+2}} + \left(\frac{y_n}{H^2} \right)^{\frac{2}{N+2}} \frac{R_0^2 z_n^{1+\kappa}}{H^2} \right\}$$

and dividing by R_0^{N+2}

$$Y_{n+1} \leq \frac{\tilde{C}2^{5n}}{H} \left\{ Y_n^{1 + \frac{2}{N+2}} + Y_n^{\frac{2}{N+2}} Z_n^{1+\kappa} \frac{R_0^{N\kappa}}{H^2} \right\}$$

Now if $H^2 < R_0^{NK}$, $\frac{R_0^{NK}}{H^2} < 1$ and [I] is proved.

To prove [II] observe that

$$z_{n+1}(k_{n} - k_{n+1})^{2} = (k_{n} - k_{n+1})^{2} \left\{ \int_{t_{0}-R_{n+1}}^{t_{0}} [meas A_{\mathbf{k}_{n+1},R_{n+1}}^{-1}(\tau)]^{\frac{r}{q}} d\tau \right\}^{\frac{2}{r}} \le$$

$$\leq ||(\mathbf{v} - k_{n})^{-\zeta_{n}}||_{q,r,\overline{Q}_{n}}^{2} \le \text{ by the embedding lemma } 2.4 \le$$

$$\leq \beta^{2} |(\mathbf{v} - k_{n})^{-\zeta_{n}}|_{\mathbf{v}_{0}^{1},\mathbf{v}_{0}^{1}(\overline{Q}_{n})}^{2}.$$

The last term is estimated in (3.7) so that [II] follow at once.

Proof of lemma 3.1 concluded:

By lemma 5.7 of [18] page 96, there exists a number 1 > 0 such that if

$$Y_1 < \lambda$$
 ; $Z_1 < \lambda^{\frac{1}{1+\kappa}}$,

then the recursion inequalities [I] - [II] imply that Y_n , $Z_n = 0$ as $n = \infty$. From [18] setting

$$d = \min\{\frac{2}{N+2}; \frac{\kappa}{1+\kappa}\}$$
,

the number 1 is given by

$$V = \min \left\{ \left(\frac{H}{2\tilde{C}} \right)^{\frac{N+2}{2}} - \frac{5(N+2)}{2\tilde{d}} ; \left(\frac{H}{2\tilde{C}} \right)^{\frac{1+\kappa}{\kappa}} - \frac{5}{\kappa \tilde{d}} \right\}.$$

Now since $\kappa \approx \frac{2\kappa_1}{N}$, $\kappa_1 = (0,1)$ and $\frac{H}{C} < 1$, above gives

$$d = \frac{1}{1+\kappa}; \quad \frac{1}{2} \cdot 0 = \left(\frac{1}{2^{s+2} \cdot \hat{c}}\right)^{\frac{1+\kappa}{\kappa}} \min \left\{\frac{-\frac{5(N+2)}{2d}}{2}, \frac{-\frac{5}{\kappa}d}{2^{\kappa}d}\right\} \cdot m^{\frac{1+\kappa}{\kappa}}.$$

Where (1) has been used. Set:

$$\sum_{n = 0}^{\infty} = \left(\frac{1}{2^{s+2} \tilde{c}}\right)^{\frac{1+\kappa}{\kappa}} \quad \min \left\{\frac{-5(N+2)}{2^{d}}, 2^{\kappa d}\right\}$$

and notice that c_0 depends only upon the data and not upon R_0 nor ω .

The lemma follows if
$$Y_1 \leq \kappa_n c_0$$
 , i.e. if

$$Y_1 = \frac{Y_1}{H^2 R_0^{N+2}} + \frac{1}{H^2 R_0^{N+2}} \int_{\mathbb{R}_0}^{\infty} \int |(v - (... + \frac{2}{2^s}))^{-1}|^2 dx d\tau \le$$

$$\frac{1}{\frac{1}{N+2}} \operatorname{meas} Q_{R_0}^{-} (n^{-} + \frac{1}{2^{S}}) = \frac{1}{P_0^{N+2}} c_0^{-} \omega^{-\frac{2\kappa_1 + N}{2\kappa_1}}, N_0^{N+2} = 0$$

from now on, for simplicity of notation we set

$$h = \frac{2 \cdot 1 + N}{2 \cdot 1}$$
, $h = \frac{1}{2 \cdot 1}$

and remark that b depends only upon the data and not upon $\ensuremath{\omega}$ nor $\ensuremath{R_0}$.

Remark: By selecting in inequalities (2.7) the constant γ large enough, we see that in [I] - [II] the constant \tilde{C} can be made as large as we please so that without loss of generality we might assume that

$$c_0 \omega^b \leq \frac{1}{2}$$
.

Suppose now that the assumption of lemma 3.1 fails; then since

$$\mu^{+} - \frac{\omega}{2^{S}} > \left| \frac{\omega}{2^{S}} + \mu^{-} \right|$$

$$\text{meas } Q_{R_{0}}^{+} (\mu^{+} - \frac{\omega}{2^{S}}) \leq \kappa_{N} R_{0}^{N+2} - \theta_{0} \kappa_{N} R_{0}^{N+2} .$$

$$= (1 - \theta_{0}) \kappa_{N} R_{0}^{N+2} .$$

Lemma 3.2: Suppose that $k \in \mathbb{R}^+$ and that

meas
$$\ensuremath{\text{Q}}_{R_{\ensuremath{0}}^{+}}^{+}(\ensuremath{k}) \; \buildrel \leq \; (1 \; - \; \ensuremath{\theta_{\ensuremath{0}}}) \ensuremath{\kappa}_{N} \; \ensuremath{R_{\ensuremath{0}}^{N+2}} \ensuremath{\ensuremath{0}}_{0} \ensuremath{\ensuremath{0}}_{N} \; \ensuremath{\ensuremath{\eta}}_{N} \; \ensuremath{\ensuremath{\eta$$

then for every $\alpha \in (0, \theta_0)$, there exists

$$\tau \in [t_0 - R_0^2, t_0 - \alpha R_0^2]$$
,

such that

meas
$$A_{k,R_0}^+(\tau) \leq \frac{1-\theta_0}{1-\alpha} \kappa_N R_0^N$$
.

Proof of lemma 3.2: If not, for all $\tau \in [t_0 - R_0^2, t_0 - \alpha R_0^2]$

meas
$$A_{k,R_0}^+(\tau) > \frac{1-\theta_0}{1-\alpha} \kappa_N R_0^N$$
 and

meas
$$Q_{R_0}^+(k) \ge \int_{t_0-R_0^2}^{t_0-\alpha R_0^2} meas A_{k,R_0}^+(\tau) d\tau >$$

$$> (1 - \theta_0) \kappa_N R_0^{N+2}$$
.

We will choose

$$\alpha = \frac{\theta_0}{2} = \frac{c_0 \omega}{2} .$$

and use the previous lemma for the levels $k=\frac{\mu}{2}^+-\frac{\omega}{2^p}$, $\forall p\,\geq\,s\,.$

Lemma 3.3: Let $\alpha = \frac{\theta_0}{2} = \frac{c_0 \omega^b}{2}$ and consider the cylinder

$$Q_{R_0}^{\alpha} = \{ |\mathbf{x} - \mathbf{x}_0| < R_0 \} \times [t_0 - aR_0^2, t_0]$$
.

There exists $p_0 \in \mathbb{N}$ dependent upon a (and hence 1) such that if

$$\frac{1}{p_0} \geq R_0^{N\kappa/2} \quad , \quad \text{then}$$

meas
$$A_{\mu^{+}}^{+} - \frac{\omega}{\frac{p_{0}}{2}}$$
 , R_{0}^{-} (t) < $[1 - \left(\frac{\theta_{0}}{2}\right)^{2}] \kappa_{N}^{-} R_{0}^{N}$,

for all $t \in [t_0 - \pi R_0^2, t_0]$.

Proof of lemma 3.3: Consider lemma 2.2 applied to the function

$$(x,t) \rightarrow (v - (\mu^+ - \frac{\omega}{2^s}))^+ (x,t)$$
, in the cylinder

$$Q_{R_0}^{\tau} = \{ |x - x_0| < R_0 \} \times [\tau, t_0] ,$$

for $\eta = \frac{\omega}{2^p}$, $p \ge s + 2$. Here $t_0 - R_0^2 \le \tau \le t_0 - \alpha R_0^2$ is the number claimed in lemma 3.2. Notice that

ess sup
$$(v - (\mu^{+} - \frac{\omega}{2^{s}}))^{+} \leq \frac{\omega}{2^{s}}$$
,

therefore for all $t \in [\tau, t_0]$ lemma 2.2 gives

(3.8)
$$\int_{B(R_0 - \sigma_1 R_0)} \ln^{+2} \left[\frac{\frac{\omega}{2^s}}{\frac{\omega}{2^s} - (v - (\mu^+ - \frac{\omega}{2^s}))^+ + \frac{\omega}{2^p}} \right] (x,t) dx \le$$

$$<\int_{B(R_0)}^{2\pi} \ln^{+2} \left[\frac{\frac{\omega}{2^s}}{\frac{\omega}{2^s} - (v - (u^+ - \frac{\omega}{2^s}))^+ + \frac{\omega}{2^p}} \right] (x, \tau) dx +$$

$$+ \frac{c}{\sigma_1^2} \left(1 + \ln \frac{\frac{\omega}{2^s}}{\frac{\omega}{2^p}} \right) \left(1 + \frac{R_0^{N\kappa}}{\left(\frac{\omega}{2^p}\right)^2} \right) \kappa_N R_0^N .$$

Let $p_0 > p$ to be selected, then if $\frac{\omega}{p_0} \ge \frac{\frac{N\kappa}{2}}{2}$, the last term is

majorized by

$$\frac{4C}{\sigma_1^2} \, \ln \, 2^{p-s} \, \stackrel{\searrow}{\sim}_N \, \stackrel{N}{R_0} \quad , \quad p \stackrel{\searrow}{\sim} s \, + \, 2 \quad .$$

We estimate the remaining terms in (3.8) as follows.

$$\int_{B(R_{0})} 2n^{+2} \left[\frac{\frac{\omega}{2^{s}}}{\frac{\omega}{2^{s}} - (v - (i^{2} - \frac{\omega}{2^{s}}))^{+} + \frac{\omega}{2^{p}}} \right] (x, t) dx \le$$

$$\leq$$
 [ln 2^{p-s}]² meas A⁺ $u^+ - \frac{\omega}{2^s}$, R₀ (\tau) \leq by lemma 3.2 \leq

$$\leq \left[(2n \ 2^{p-s})^2 \left(\frac{1 - \theta_0}{1 - \alpha} \right) \kappa_N^{N} R_0^N \right].$$

For the left hand side we have

$$\int_{B(R_0^{-\sigma_1}R_0^{\sigma_0})}^{2n^{+2}} \left[\frac{\frac{\omega}{2^s}}{\frac{\omega}{2^s} - (v - (\mu^+ - \frac{\omega}{2^s}))^+ + \frac{\omega}{2^p}} \right] (x,t) dx \ge$$

$$\geq \int_{B(R_0 - \sigma_1 R_0) \cap \{v > \mu^+ - \frac{\omega}{2^p}\}} \left[\frac{\frac{\omega}{2^s}}{\frac{\omega}{2^s} - (v - (\mu^+ - \frac{\omega}{2^s}))^+ + \frac{\omega}{2^p}} \right] (x, t) dx \geq$$

$$\geq \left[2n \left(\frac{\frac{\omega}{2^{s}}}{2 \frac{\omega}{2^{p}}} \right) \right]^{2} \operatorname{meas} A^{+}_{\mu^{+} - \frac{\omega}{2^{p}}, R_{0}^{-\sigma_{1}R_{0}}} (t) .$$

These estimates in (3.8) give the inequality

(3.9) meas
$$\frac{1}{p^{+}} + \frac{\pi}{2^{p}} R_{0} - \sigma_{1} R_{0}$$
 (t) $\leq \left(\frac{2n 2^{p-s}}{2n 2^{p-s-1}}\right)^{2} \left(\frac{1 - R_{0}}{1 - R_{0}}\right) R_{0} R_{0}^{N} + \frac{\pi}{2^{p-s}}$

$$\begin{split} & + \frac{4C}{\sigma_1^2} \frac{\ln 2^{p-s}}{(\ln 2^{p-s-1})^2} \, \kappa_N \, \, R_0^N = \\ & = \left(\frac{p-s}{p-s-1} \right)^2 \, \left(\frac{1-\theta_0}{1-\alpha} \right) \, \kappa_N \, \, R_0^N + \frac{4C}{\sigma_1^2 \, \ln 2} \, \frac{p-s}{(p-s-1)^2} \, \kappa_N \, \, R_0^N \quad . \end{split}$$

Now

meas
$$A^{+}_{\mu^{+}} - \frac{\omega}{2^{p}}, R_{0}$$
 (t) $\leq \text{meas } A^{+}_{\mu^{+}} - \frac{\omega}{2^{p}}, R_{0}^{-\sigma_{1}} R_{0}$ (t) +

+ meas
$$[B(R_0) \setminus B(R_0^{-\sigma_1}R_0)] \leq \max_{\mu^+} A^+ + \frac{\omega}{2^p} R_0^{-\sigma_1}R_0$$
 (t) +

$$+ N\sigma_1 \kappa_N R_0^N$$
,

therefore by virtue of (3.9)

$$\max_{\mu^{+}} A_{\frac{\omega}{2^{p}}, R_{0}}^{+} \text{(t)} \leq \left\{ \left(\frac{p-s}{p-s-1} \right)^{2} \left(\frac{1-\theta_{0}}{1-\alpha} \right) + \frac{4C}{\sigma_{1}^{2} \ln 2} \frac{p-s}{(p-s-1)^{2}} + N\sigma_{1} \right\} \kappa_{N} R_{0}^{N} .$$

This inequality holds for all $\sigma_1 \in (0,1)$, all p > s + 2 and all $t \in [\tau,t_0]$.

Select $\sigma_1 = \frac{3}{8} \frac{\theta_0^2}{N}$, and p_0 so large that

$$\frac{4C}{\sigma_1^2 \ln 2} \frac{p_0 - s}{(p_0 - s - 1)^2} < \frac{3\theta_0^2}{8} , \text{ and}$$

$$\left(\frac{p_0 - s}{p_0 - s - 1}\right)^2 \le (1 - \alpha)(1 + \alpha_0) ,$$

to obtain

meas
$$A_{\mu}^{+} = \frac{\omega}{2^{P_0}}, R_0$$
 (t) $\leq \left[1 - \left(\frac{\theta_0}{2}\right)^2\right] \kappa_N R_0^N$.
$$= \left[1 - \left(\frac{e_0 \omega}{2}\right)^2\right] \kappa_N R_0^N.$$

This proves the lemma.

Remark: It is easily seen that a suitable choice of p_0 is

(3.10)
$$p_0 = s + 3 + \left[\frac{c_1}{(c_0^b)^6} \right],$$

where [a] denotes the largest integer contained in a, and

$$c_1 = \frac{2^8 N^2 C}{\ln 2}$$

Notice that C_1 depends only upon the data and not upon ω nor R_0 .

Remark: Since for $q \ge p_0$, $A^+_{\mu} + \frac{\omega}{2^q}$, R^-_0 (t) $\in A^+_{\mu} + \frac{\omega}{2^n}$, R^-_0 (t), we have that

meas
$$A_{\mu}^{+} = \frac{1}{2^{q}} R_{0}^{R}$$
 (t) $\leq \left[1 - \left(\frac{\theta_{0}}{2}\right)^{2}\right] K_{N} R_{0}^{N}$, $\forall q \geq p_{0}$,

and for all $t \in [t_0 - \alpha R_0^2, t_0]$.

The subsequent arguments will be carried over the cylinder $Q_{R_0}^\alpha$. For k > 0 we also denote

$$Q_{R_{0}}^{(k)}(k) = \{(x,t) \in Q_{R_{0}}^{\alpha}|v(x,t) > k\}$$
.

<u>Lemma 3.4</u>: For every $\theta_1 > 0$, there exists $q_0 \in \mathbb{N}$, $q_0 > p_0$ such that if

$$\frac{\omega}{q_0} > R_0^{\frac{N\kappa}{2}}$$
 , then

meas
$$Q_{R_0}^{\alpha}(\mu^+ - \frac{\omega}{q_0}) \leq \theta_1 \kappa_N R_0^{N+2}$$
.

Proof of lemma 3.4: Lemma 3.3 and the remarks following it imply that

$$\operatorname{meas}\{B(R_0) \setminus A^+ \atop \mu^+ - \frac{\omega}{2^q}, R_0^{(t)}\} \ge \left(\frac{\theta_0}{2}\right)^2 \kappa_N R_0^N , \quad \forall q \ge P_0$$

for all $t \in [t_0 - \alpha R_0^2, t_0]$.

Apply inequality (2.9) to the function $x \to v(x,t)$ in the ball $B(R_0) \times \{t\} \ \text{for the levels}$

$$\ell = \mu^{+} - \frac{\omega}{2^{q+1}}$$
, $k = \mu^{+} - \frac{\omega}{2^{q}}$, $q_0 > q > p_0$,

where q_0 has to be chosen. If we do this for all $t \in [t_0 - \alpha R_0^2, t_0]$ we obtain

$$\int_{A_{k,R_0}^+(t)\setminus A_{\ell,R_0}^+(t)} |\nabla_{\mathbf{x}} \mathbf{v}| d\mathbf{x} \leq \frac{4DR_0}{\kappa_N^{\theta_0}} \int_{A_{k,R_0}^+(t)\setminus A_{\ell,R_0}^+(t)} |\nabla_{\mathbf{x}} \mathbf{v}| d\mathbf{x} .$$

Integrate both the sides of this inequality over $[t_0 - aR_0^2, t_0]$, square and use Hölder's inequality on the right side, to obtain

$$(3.11) \qquad \left(\frac{\omega}{2^{q+1}}\right)^{2} \left[\text{ meas } Q_{R_{0}}^{\alpha}\left(\mu^{+} - \frac{\omega}{2^{q+1}}\right)\right]^{2} \leq \left[\frac{4D}{e_{0}^{2}\kappa_{N}}\right]^{2} R_{0}^{2} .$$

$$\left[\int_{t_{0}^{-1}A_{0}^{2}}^{t_{0}} A_{k,R_{0}}^{+}(\tau) A_{k,R_{0}}^{+}(\tau)\right]^{2} dxd\tau \left[\int_{t_{0}^{-1}A_{0}^{2}}^{t_{0}} [\text{meas } A_{k,R_{0}}^{+}(\tau) A_{k,R_{0}}^{+}(\tau)]d\tau\right]$$

$$\leq \left[\frac{4D}{e_{0}^{2}\kappa_{N}}\right]^{2} R_{0}^{2} \left[\left(\nu - \left(\mu^{+} - \frac{\omega}{2^{q}}\right)\right)^{+}\right]^{2} v_{2}^{1}, o\left(Q_{R_{0}}^{2}\right)$$

$$\left[\int_{t_{0}^{-1}A_{0}^{2}}^{t_{0}} [\text{meas } A_{k,R_{0}}^{+}(\tau) A_{k,R_{0}}^{+}(\tau)]d\tau\right] .$$

In order to estimate the $v_2^{1,0}(Q_{R_0})$ -norm of $(v-(\mu^+-\frac{\omega}{2^q}))^+$ we use inequalities (2.7) applied to the pair of cylinders Q_{R_0} , Q_{2R_0} . Notice that in this connection ess sup $(v-(\mu^+-\frac{\omega}{2^q}))^+\leq \frac{\omega}{2^q}<\delta$ and that Q_{2R_0}

$$(\sigma_1 R_0)^{-2} = 4R_0^{-2} ; (\sigma_2 R_0^2)^{-1} = \frac{4}{3} R_0^{-2}$$
.

Moreover observe that since $\mu^{+} \geq |\mu^{-}|$, we have

$$\mu^+ - \frac{\omega}{2^q} > 0 \quad ,$$

so that $\phi_a^+(\hat{x} - \frac{\omega_0}{2^q}, t_0 - 4R_0^2, t) = 0$, $t \in [t_0 - 4R_0^2, t_0]$. Inequalities (2.7) now give

$$\| (\mathbf{v} - (\mathbf{b}^{+} - \frac{1}{2^{\mathbf{q}}}))^{+} \|_{\mathbf{V}_{2}^{1,0}(Q_{\mathbf{R}_{0}})}^{2} \leq \gamma (4 + \frac{4}{3}) R_{0}^{-2} \| (\mathbf{v} - (\mathbf{b}^{+} - \frac{1}{2^{\mathbf{q}}}))^{+} \|_{2,Q_{2\mathbf{R}_{0}}}^{2} +$$

$$+ \gamma \left\{ \begin{array}{c} t_0 \\ \int_{t_0 - 4R_0^2} \left[\max_{\mu} A^+ - \frac{\omega}{2^q}, 2R_0 \right]^{\frac{r}{q}} d\tau \end{array} \right\} \stackrel{\frac{2}{r}}{(1+\kappa)} \leq$$

$$\leq \gamma \frac{2^{N+6}}{3} \left(\frac{\omega}{2^{q}}\right)^{2} \kappa_{N} R_{0}^{N} + \gamma 2^{N(1+\kappa)} \kappa_{N}^{\frac{2}{q}} (1+\kappa)^{-1} R_{0}^{N\kappa} \kappa_{N} R_{0}^{N}$$
,

where (2.4) has been used. By assumption

$$R_0^{N\kappa} \leq \left(\frac{\omega}{\frac{q}{2}}\right)^2 < \left(\frac{\omega}{2^q}\right)^2$$
,

so that there exists a constant C_2 depending only upon the dimension N and the data, such that

$$|(v - (\mu^+ - \frac{\omega}{2^q}))^+|_{V_2^1, O(Q_{R_0})}^2 \le C_2 \left(\frac{\omega}{2^q}\right)^2 \kappa_N^N R_0^N$$
.

Carrying this in (3.11) and dividing by $\left(\frac{\omega}{2^{q+1}}\right)^2$, gives

(3.12)
$$\left[\text{meas } Q_{R_0}^{\alpha} \left(\mu^+ - \frac{\omega}{2^{q+1}}\right)\right]^2 \leq 4C_2 \left[\frac{4D}{\kappa_N}\right]^2 \frac{1}{\theta_0^4} \kappa_N R_0^{N+2}$$
.

$$\begin{bmatrix} t_0 \\ \int_{t_0-\alpha R_0^2} [\text{meas } A_{k,R_0}^+(\tau) \setminus A_{\ell,R_0}^+(\tau)] d\tau \end{bmatrix}.$$

We add inequalities (3.12) with respect to q, from p_0 to q_0-1 and obtain

$$(q_0 - p_0) [\text{meas } Q_{R_0}(u^+ - \frac{\omega}{q_0})]^2 \le 4C_2 [\frac{4D}{\kappa_N}]^2 \frac{1}{\theta_0^4} \kappa_N R_0^{N+2}$$
.

$$\sum_{q=p_0}^{q_0-1} \int_{t_0-\alpha R_0^2}^{t_0} [meas A^+_{\mu^+} - \frac{\omega}{2^q}, R_0^{(\tau)}] A^+_{\mu^+} - \frac{\omega}{2^{q+1}}, R_0^{(\tau)}] d\tau \le$$

$$\leq 4C_2 \left[\frac{4D}{\kappa_N}\right]^2 \frac{1}{\theta_0^4} \alpha (\kappa_N R_0^{N+2})^2$$
.

Now recall that $\alpha = \frac{\theta_0}{2}$ set

$$C_3 = 2C_2 \left[\frac{4D}{\kappa_N}\right]^2$$

and observe that C_3 depends only upon the data and the dimension N.

Dividing the inequality above by q_0 - p_0 , to prove the lemma we have only to choose q_0 so large that

$$\frac{1}{q_0 - p_0} \frac{c_3}{\theta_0^3} \le \theta_1^2 \quad .$$

We will select

(3.13)
$$q_0 = p_0 + 1 + \begin{bmatrix} \frac{c_3}{3} \\ \frac{3}{0} \\ 0 \end{bmatrix}^2$$

Remark: The proof of lemma 3.4 is an adaptation of a similar result of [18], namely lemma 7.2 page 114.

Consider now the pair of cylinders $Q_{R_0}^{\alpha}$ and

$$\frac{2_{R_0}^{+}}{\frac{9}{2}} + |\mathbf{x} - \mathbf{x}_0| + \frac{R_0}{2} + |\mathbf{ft}_0| - |\mathbf{x}| \frac{R_0^2}{4}, |\mathbf{t}_0|.$$

For them we have the following result

Lemma 3.5: There exists a number $\theta_1 > 0$ depending upon α , N and the data, such that if

meas
$$Q_{R_0}^{\alpha} (\mu^+ - \frac{\omega}{q_0}) < \theta_1 \kappa_N R_0^{N+2}$$
 ,

then either

(i)
$$H = \text{ess sup } (v - (\mu^+ - \frac{\omega}{q_0}))^+ \le \frac{N\kappa}{2}$$
, or $Q_{R_0}^{\alpha}$

(ii) meas
$$Q_{\frac{R}{2}}^{\alpha} (\mu^{+} - \frac{\omega}{q_{0}} + \frac{1}{2} H) = 0$$
.

<u>Proof of lemma 3.5</u>: The proof is very similar to the proof of lemma 3.1. We reproduce the main steps mainly to trace the dependence of q_0 on q_0 (and hence on ω). Let R_n , \overline{R}_n be defined as before, and consider the cylinders

$$Q_{R_n}^{\alpha} = \{ |\mathbf{x} - \mathbf{x}_0| < R_n^{\gamma} \times [t_0 - \alpha R_n^2, t_0] \}$$

$$\bar{Q}_{n}^{\alpha} \equiv \{|\mathbf{x} - \mathbf{x}_{0}| < \bar{\mathbf{R}}_{n}\} \times [\mathbf{t}_{0} - \alpha \mathbf{R}_{n}^{2}, \mathbf{t}_{0}]$$

$$\overline{\overline{Q}}_{n}^{\alpha} = \{ |\mathbf{x} - \mathbf{x}_{0}| < \overline{R}_{n} \} \times [\mathbf{t}_{0} - \alpha R_{n+1}^{2}, \mathbf{t}_{0}] ,$$

which satisfy the inclusions

The second second

$$Q_{R_{n+1}}^{\alpha} \subset \overline{Q}_{n}^{\alpha} \subset \overline{Q}_{n}^{\alpha} \subset Q_{R_{n}}^{\alpha}$$
.

We use inequalities (2.7) over $\overline{\mathbb{Q}}_n^{x}$ and \mathbb{Q}_{R}^{x} , for the functions $(x,t)+(v-k_n)^+$ where

$$k_n = (k_1 + \frac{1}{2}H) - \frac{1}{2^n}H$$
 , $n = 1, 2, ...$

$$k_1 = \frac{1}{2} - \frac{\omega}{q_0}$$
.

Since $k_n \ge k_1 > 0$ in (2.7) we have

$$\phi_{a}^{+}(k_{n},t_{0}-\alpha R_{n}^{2},t_{0})=0$$
 .

Note that in this case $(\mathfrak{I}_1 \mathbb{R}_n)^{-2} = \mathbb{R}_0^{-2} 2^{2(n+4)}$ and $(\mathfrak{I}_2 \mathfrak{a} \mathbb{R}_n^2)^{-1} = \mathbb{R}_0^{-2} \mathfrak{a}^{-1} 2^{n+3}$. We have to show that the numbers

$$Y_n = \frac{1}{H^2 R_0^{N+2}} Y_n = \frac{1}{H^2 R_0^{N+2}} \int_{Q_{R_n}}^{\alpha} \int (v - k_n)^{+2} dx d\tau$$

$$Z_{n} = \frac{z_{n}}{R_{0}^{N}} = \frac{1}{R_{0}^{N}} \left\{ \int_{R_{0}^{-\alpha R_{n}}}^{t_{0}} \left[\max_{k_{n}, R_{n}} A_{n}^{+}(\tau) \right]^{\frac{r}{q}} d\tau \right\} \frac{2}{r},$$

tend to zero as n $^{+}$ $^{\infty}$. Proceeding exactly as in lemma 3.1 we see that Y n and Z satisfy the recursion inequalities

$$Y_{n+1} \le \frac{\tilde{c}2^{4n}}{L} [Y_n^{1+\frac{2}{N+2}} + Y_n^{\frac{2}{N+2}} z_n^{1+\kappa}]$$

$$z_{n+1} \le \frac{\hat{c}2^{4n}}{1} [Y_n + Z_n^{1+\kappa}]$$
.

The lemma follows if

$$Y_1 > F_0^{-(N+2)} \text{ meas } P_{K_0}^{-(k_1)} \le 1 + N =$$

$$=\left(\frac{1}{2\tilde{c}}\right)^{\frac{1+\kappa}{\kappa}}\frac{1+\kappa}{\alpha}\min\{2^{-4}\frac{N+2}{2\tilde{d}},2^{-4}\frac{1}{\kappa\tilde{d}}\},$$

i.e. if

$$\theta_1 = C_4 \alpha^b$$
, where $C_4 = \frac{1}{\kappa_N} (2\tilde{C})^{-\frac{1+\kappa}{\kappa}} \min\{2^{-4\frac{N+2}{2d}}, 2^{-4\frac{1}{\kappa d}}\}$.

Here \tilde{C} and d are as in lemma 3.1.

From (3.13) and the remarks above it follows that the conclusion of lemma 3.5 holds true if we choose

$$q_0 \ge p_0 + 1 + \left[\frac{c_3}{(c_4 a^b)^2 \theta_0^3} \right]$$
.

We recall that $\theta_0 = c_0 \omega^b$, $\alpha = \frac{\theta_0}{2}$, therefore taking in account (3.10) we might select

$$q_0 = s + 4 + \begin{bmatrix} c_1 + \frac{2^{2b}c_3}{c_4^2} \end{bmatrix} \frac{1}{(c_0\omega^b)^{\max\{6,2b+3\}}}$$

Set

$$A = \begin{cases} c_1 + \frac{2^{2b}c_3}{c_4^2} \end{cases} \frac{1}{c_0^{\max\{6,2b+3\}}}$$

and

$$a = b \max\{6 : 2b + 3\}$$

so that

(3.14)
$$q_0 = s + 4 + \left[\frac{A}{\omega^a}\right].$$

We remark that s, A, a depend only upon the data and not upon $\,\omega$ nor the size of $\,E_{0}^{}$.

Front of the proposition: Suppose that

(3.15)
$$\frac{1}{2^{s}} \cdot \frac{q_{0}+1}{2} = \frac{1}{s+4+A/\omega^{a}} \cdot (2F_{0})^{\frac{N}{2}}.$$

Chviously either

1. meas
$$Q_{R_0}^- (1^- + \frac{\omega}{2^s}) \le c_0^- \frac{\frac{2^s_1}{N+2^s_1}}{s_N^{N+2}}$$

or

2. meas
$$Q_{R_0}^- (N_0^- + \frac{1}{2^s}) > C_0^- \frac{\frac{2^k 1}{N+2^k 1}}{k_N^2 R_0^{N+2}}$$
.

Case 1: By lemma 3.1 either

1.a. ess sup
$$(v - (L^2 + \frac{\omega}{2^s}))^{-\frac{N^k}{2^k}}$$
, or Q_{R_0}

1.b. meas
$$Q_{\frac{R}{2}}(1 + \frac{\omega}{2^s} - \frac{1}{2}H) = 0$$
.

If l.a occurs then

- ess inf
$$v \le -\mu - \frac{\omega}{2^s} + \frac{N\kappa}{2} \le by$$
 (3.15) $\sim \Omega_{R_0}$

$$<-1^{-}-\frac{1}{2^{s+5+A/a}}$$
.

Adding each sup voon the left hand side and u^{+} on the right hand side we $\Omega_{P_{0}}$

obtain

PSS OSC
$$V < GSS$$
 OSC $V - \frac{\omega}{2^{S+5+A/\omega^a}}$

$$\leq \omega \left(1 - \frac{1}{2^{S+5+A/\omega^a}}\right).$$

If 1.b occurs then

- ess inf
$$v \le - \nu^{-} - \frac{\omega}{2^{s}} + \frac{1}{2} \left(- \nu^{-} + \frac{\omega}{2^{s}} + \nu^{-} \right)$$

$$= -1.^{-} - \frac{3}{2^{s+1}}$$

i.e.

ess osc
$$v \le \omega(1 - \frac{1}{2^{s+5+A/\omega^a}})$$
.

Case 2: By lemmas 3.2 - 3.4, in view of (3.15), the assumptions of lemma 3.5 are verified. It gives the following alternative. Either

2.a. ess sup
$$v \leq \frac{1}{2} + \frac{\frac{N\kappa}{2}}{\frac{q}{2}} \leq \frac{N\kappa}{2} \leq \frac{N\kappa}{2}$$

or

2.b. ess sup
$$v = u^{+} - \frac{1}{q_{0}} + \frac{1}{q_{0}+1} = u^{+} - \frac{1}{q_{0}+1}$$
.
$$\frac{2^{1}}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} =$$

Hence in either case

ess osc
$$v \leq \omega (1 - \frac{1}{2^{s+5+A/\omega}})$$
.
$$\frac{2^{\frac{1}{R_0}}}{2}$$

Now to determine R_{\star} notice that by virtue of (3.15)

$$a\left(\frac{R_0}{2}\right)^2 = \frac{1}{2 \cdot 4^2} c_0 \omega^b (2R_0)^2 >$$

$$> 2^{\text{sb-5}} c_0(2R_0)^{2+\frac{N\kappa b}{2}}$$
.

Setting $t_{\star} = (2^{\text{sb-5}}c_0)^{\frac{1}{2}} + 1$, and

$$h = 1 + \frac{NKb}{4} ,$$

we have

$$\sqrt{2} \frac{R_0}{2} \ge 5*(2R_0)^h = R_*$$

so that $Q_{R_0}^2 = Q_{R_*} = \{ [x + x_0] < R_* \} \times [t_0 - R_*^2, t_0].$

It follows that

where $s_0 = s + 5$. Finally if (3.3) is false then (3.16) follows at once. The proposition is proved.

4. Proof of Theorem 1:

We will prove the theorem by exploiting the results of the previous section. Proposition 3.1 is valid for any number $\,\omega\,$ satisfying

ess osc v
$$\leq \omega \leq 2M$$
 . Q_{2R_0}

We stress the fact that the constants ξ_{\star} , a, b, A, h in proposition 3.1 do not depend upon R_0 nor ω . Let $(\mathbf{x}_0,\mathbf{t}_0)\in\Omega_{\mathbf{T}}$ be fixed and select $\omega=2\mathrm{M}\geq\cos v$. Let $0\leq R_0\leq\frac{1}{2}$ be so small that

$$Q_{2R_0} \subset \Omega_T$$

(4.1)
$$(2R_0)^{\frac{N\kappa}{2}} \leq \frac{2M}{s_0^{+A/(2M)^a}} .$$

Define two sequences of positive real numbers $\{{\tt R}_n\}$ and $\{{\tt M}_n\}$ as follows

$$p_1 = 2p_2$$
; $p_{n+1} = p(p_n)^h$ $n = 2, 3, ...$

where $\gamma = \min\{\zeta_*, 4^{-\frac{2a}{11\zeta_*}}\}$, and

$$x_1 = 2x$$
, $x_{n+1} = x_n (1 - \frac{1}{2^{s_0 + A/M_n^a}})$

Lemma 4.1: $\{(n_n + 1, n_n) + 1, (n_n + 1) \text{ and for all } n \in \mathbb{N}\}$

$$\lim_{n \to \infty} \frac{1}{2} \operatorname{dist}_{n} = \sum_{i=1}^{n} \frac{1}{n} \operatorname{dist}_{n} = \sum_{i=1}^{n} \operatorname{dist}_{n} = \sum_{i=1}^{n} \frac{1}{n} \operatorname{dist}_{n} = \sum_{i=1}^{n} \frac{$$

Proof of Lemma 4.1: If $M_n \ge M_{n+1} \ge ... \ge M_0 \ge 0$, then for all $n \in \mathbb{N}$

$$M_{n+1} \leq M_n \varepsilon$$
; $\varepsilon = (1 - \frac{1}{2^{s_0 + A/M_0^a}}) \leq 1$.

Phorefore

$$M_n \leq M_0 \varepsilon^n$$
 $n = 1, 2, ...$

which implies that M $_n$ V 0 as n $^{\to}$ $^{\infty}.$ A contradiction. The statement about $\{E_n\}$ is trivial.

In view of 4.1, Proposition 3.1 implies that

ess osc v
$$\leq M_2$$
 . Q_{R_2}

Moreover

$$R_{2}^{\frac{N\kappa}{2}} = \xi^{\frac{N\kappa}{2}} \left(R_{1}^{\frac{N\kappa}{2}} \right)^{h} \leq 4^{-a} \left(\frac{M_{1}}{2^{s_{0} + A/M_{1}^{a}}} \right)^{h}$$

For simplicity of notation set

$$s_0^{+A/x^a}, x > 0.$$

Fig., using the definition of $M_{\gamma\gamma}$

$$\frac{N_{1}}{K_{2}^{2}} < 4^{-a} \left(\frac{M_{2}}{\sigma(M_{1})(1 - \sigma(M_{1})^{-1})} \right)^{h} \le 4^{-a} \left(\frac{2M_{2}}{\sigma(M_{1})} \right)^{h} =$$

$$= \frac{M_2}{\sigma(M_2)} \frac{\sigma(M_2)}{\sigma(M_1)} 4^{-a} 2 \left(\frac{2M_2}{\sigma(M_1)}\right) \frac{N_b b}{4}.$$

Limithful the definitions of b, s it is immediate to see that $2(\frac{2M_2}{c(M_1)})^{\frac{NKD}{4}} > 1$.

Now for all $n \in \mathbb{N}$ it is easy to check that

$$\frac{\sigma(M_{n+1})}{\sigma(M_n)} \leq 4^{a} ,$$

Hence

$$R_2^{\frac{N\kappa}{2}} \leq \frac{M_2}{2^{s_0+A/M_2^a}}.$$

We have shown that the two inequalities

$$\begin{array}{c}
\operatorname{osc} v \leq M_{\mathbf{I}} \\
Q_{R_{\mathbf{I}}}
\end{array}$$

$$R_{1}^{\frac{N\kappa}{2}} \leq \frac{M_{1}}{2^{s_{0}+A/M_{1}^{a}}}$$

imply the same two inequalities for $\rm R_2$ and $\rm M_2$. The same argument shows that if

$$\operatorname{ess\ osc\ } v \, \leq \, \operatorname{{\tt M}}_n$$

$$\frac{\frac{N\kappa}{2}}{R_n} \leq \frac{\frac{M}{n}}{\frac{s}{2}0^{+A/M_n^a}},$$

then the same inequalities are valid for n + 1. The lemma is proved.

As a consequence of lemma 4.1 we have that $\forall (\mathbf{x}_0, \mathbf{t}_0) \in \mathbb{Q}_{\mathbf{T}}$

ess lim
$$v(x,t)$$

 $(x,t)+(x_0,t_0)$

exists. We define the function $(x,t) \rightarrow \hat{v}(x,t)$ by setting

$$\hat{\mathbf{v}}(\mathbf{x}_0, \mathbf{t}_0) = \underset{(\mathbf{x}, \mathbf{t}) + (\mathbf{x}_0, \mathbf{t}_0)}{\text{oss lim}} \mathbf{v}(\mathbf{x}, \mathbf{t})$$
.

Lemma 4.2: The function $(\mathbf{x},\mathbf{t}) + \hat{\mathbf{v}}(\mathbf{x},\mathbf{t})$ is a continuous representative out of the equivalence class \mathbf{v} . Moreover if \mathbf{K} is a compact contained in $\mathbf{v}_{\mathbf{T}}$ there exists a nondecreasing continuous function $\mathbf{w}_{\mathbf{K}}(\cdot) \colon \mathbb{R}^+ \times \mathbb{R}^+, \mathbf{w}_{\mathbf{K}}(0) = 0$ depending upon the data and dist($\mathbf{K} : \mathbf{r}$) such that

$$|\hat{\mathbf{v}}(\mathbf{x}_1, \mathbf{t}_1) - \hat{\mathbf{v}}(\mathbf{x}_2, \mathbf{t}_2)| \le \varepsilon_K (|\mathbf{x}_1 - \mathbf{x}_2| + |\mathbf{t}_1 - \mathbf{t}_2|^{\frac{1}{2}})$$

$$\forall (\mathbf{x}_i, \mathbf{t}_i) \in K , i = 1, 2 .$$

The statement is a direct consequence of lemma 4.1 and establishes the interior regularity claimed by Theorem 1.

Remark: The continuity result is a consequence of inequalities (2.7) - (2.8) and lemma 2.2 solely.

5. Continuity up to the boundary:

Let $u\in W_2^{1,1}(\Omega_T)$ be a weak solution of (1.1) subject to certain boundary conditions. In this section we investigate under what circumstances the continuity of u can be extended to the closure of Ω_T . As remarked in the introduction, this is equivalent to prove the continuity of $(x,t) \to v(x,t)$ on the closure $\overline{\Omega}_T$.

Our study is divided in three parts. First we show the continuity of v at time t = 0. Then we investigate the regularity on the lateral part S_T of the parabolic boundary of Ω_T , for the cases of variational (Neumann) boundary conditions, and homogeneous Dirichlet boundary data. The method of proof in all three cases is similar to the one in section 3 and 4, and consists roughly speaking in constructing for every $(x_0,t_0)\in\bar{\Omega}_T$ a family of coaxial nested cylinders with same "vertex" (x_0,t_0) where the essential oscillation of v progressively decreases according to the rules imposed by the operator in (1.1) and the boundary data. Escause of the information contained in the boundary data the analysis in the present situation is in fact simpler.

We will consider cylinders of two types.

Basis cylinders:

(BC)
$$(\mathbf{x}_0, \mathbf{t}_0) \in \Omega_T$$
; $Q_R = \{|\mathbf{x} - \mathbf{x}_0| < R\} \times [\mathbf{t}_0 - R^2, \mathbf{t}_0]\}$

with
$$\{|x - x_0| \le R\} \in \Omega$$
 and $t_0 - R^2 \le 0$.

Lateral cylinders:

(LC)
$$(x_0, t_0) \in S_T$$
; $Q_R = \{ \{x - x_0 \} < R\} \times \{t_0 - R^2, t_0 \}$.

The axis of (LC) lies on ${\bf S}_{\widehat{T}}$ and both (BC) and (LC) are not contained in ${\bf G}_{\widehat{T}}$.

A!-Continuity at t = 0

Let $\mathbf{v} \in \mathbb{V}_2^{1,1}(\mathbb{T})$ satisfy identity (2.1) and in addition

$$v(x,0) = v_0(x) = \hat{s}_0^{-1}(u_0(x))$$
,

in the sense of the traces over $\mathbb S$. Let the selection $w_0(x) \in \mathcal C(u_0(x))$ be given so that (2.1) is well defined for all $t_0 \ge 0$. We assume that $x + v_0(x)$ is continuous in $\mathbb S$, with modulus of continuity $s + \omega_K(s)$ over a compact $K \subset \mathbb N$. Here $\mathbb S_K(s)$ maps $\mathbb R^+ \to \mathbb R^+$ is non-decreasing and $\mathbb S_K(0) = 0$. Our task is to prove the following theorem.

Theorem 5.1: Let K be a compact of Ω . There exists a non-decreasing, continuous function $s \to \omega_K^+(s) : \mathbb{R}^+ \to \mathbb{R}^+$, $\omega_K^-(0) = 0$ such that

$$|v(x_1,t_1) - v(x_2,t_2)| \le \omega_K(|x_1 - x_2| + |t_1 - t_2|^{\frac{1}{2}})$$

for all $(x_i, t_i) \in K' \times \{0, T\}$, i = 1, 2 , and every compact $K' \subseteq K$.

The function $\neg_K(.)$ depends only upon the data and the modulus of continuity of $|v_0(x)|$.

Clearly we have only to prove the continuity at t=0, so that it will suffice only to consider cylinders (BC).

Learn 5.1: Let $x_0 \in \Sigma$ and R > 0 so small that $B(R) \in \mathbb{R}$. Fix $0 < t_0 < R^2$ and over the cylinder (BC) Q_R consider the cut off function $x + \xi(x)$ selected as in (b). There exists a constant γ depending only upon the data such that the functions $(s,t) + (v-k)^{\frac{1}{2}}(x,t)$ satisfy the inequalities

$$\|\nabla_{\mathbf{v}}(\mathbf{v}-\mathbf{k})^{\frac{1}{2}}\|_{2,\mathbf{B}(\mathbf{R}-\gamma_{1}\mathbf{R})}^{2}(\mathbf{t}) + \|\nabla_{\mathbf{x}}(\mathbf{v}-\mathbf{k})^{\frac{1}{2}}\|_{2,\mathbf{B}(\mathbf{R}-\gamma_{1}\mathbf{R}),\gamma_{1}}^{2}0,\mathbf{t}\|_{2,\mathbf{R}}^{2}$$

$$\int_{\mathbb{R}^{N}} (v - k)^{\frac{1}{2}} \frac{2}{2 \cdot B(R)} (0) + \int_{\mathbb{R}^{N}} (c_{1}R)^{-2} + \left(\frac{2}{2} R^{2} \right)^{-1} \left[\frac{2}{2} (v - k)^{\frac{1}{2}} \frac{2}{2 \cdot Q_{R}} (c_{1}R)^{-1} \right] + C \left(\frac{2}{2} R^{2} \right)^{-1} \left[\frac{2}{2} (v - k)^{\frac{1}{2}} \frac{2}{2 \cdot Q_{R}} (c_{1}R)^{\frac{1}{2}} \right] + C \left(\frac{2}{2} R^{2} \right)^{-1} \left[\frac{2}{2} (v - k)^{\frac{1}{2}} \frac{2}{2 \cdot Q_{R}} (c_{1}R)^{\frac{1}{2}} \right] + C \left(\frac{2}{2} R^{2} \right)^{-1} \left[\frac{2}{2} (v - k)^{\frac{1}{2}} \frac{2}{2} \frac{2}{2} Q_{R} (c_{1}R)^{\frac{1}{2}} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R^{2} \right)^{\frac{1}{2}} \left[\frac{2}{2} R^{2} \right] + C \left(\frac{2}{2} R$$

$$+ \gamma \left\{ \begin{bmatrix} t & \frac{r}{q} \\ 0 & d\tau \end{bmatrix} \frac{\frac{2}{r} (1+r)}{r} + \tilde{\Phi}_{b}^{\pm} (k,0,t), \qquad \text{for all } t \in [0,t_{0}], \right\}$$

provided that $k \in \mathbb{R}$ satisfies the restriction

(5.2)
$$\operatorname{ess\ sup\ } (v-k)^{\frac{1}{2}} < \delta .$$

Here δ is the same number introduced in (2.6) and

$$\tilde{\xi}_{b}^{\pm}(k,0,t) = \pm \int_{\Omega} v(x,\tau) \chi \{v \leq 0\} (v-k)^{\pm} \zeta^{2} dx \Big|_{0}^{t}$$

$$\pm \int_{0}^{t_{0}} \int_{\Omega} v(x,\tau) \chi \{v \leq 0\} \frac{3}{4t} (v-k)^{\pm} \zeta^{2}(x) dx d\tau .$$

Inequalities (5.1) are derived in away similar to inequalities (2.7) - (2.8), the only difference being the domain of integration. In particular, the constant γ can be taken equal to the analogous constant appearing in (2.7) - (2.8).

Next we simplify (5.1) by imposing further restrictions on the levels $\ k$. Setting

$$k_1(R) = \sup_{B(R)} v_0(x)$$
, $k_2(R) = \inf_{B(R)} v_0(x)$,

for the oscillation of $x \to v_0(x)$ in B(R) we have

osc
$$v_0(x) = k_1(R) - k_2(r) + \frac{1}{2}(R)$$
,

for a compact K contained in $\mathbb R$ and containing B(R). From now on we will keep fixed the compact K, and all the subsequent arguments will be carried over balls $B(R) \subset K$.

If in (5.1) we choose $k \ge k_1(R)$, then $(v-k)^+(x,0) = 0$. Moreover if we look at v as extended over all Q_R in a way not to exceed $k_1(R)$, then $(v-k)^+$ is identically zero over that portion of Q_R not contained in Q_R . Therefore if $k \ge k_1(R)$, the domains of integration in (5.1) can be replaced by $B(R-c_1R) \times [t_0-R^2,t_0]$ on the left hand side and Q_R on the right hand side respectively. These remarks show that the function $(x,t) \to v(x,t)$ satisfies the inequalities

$$(5.1)^{+} \qquad |(v - k)^{\pm}|^{2} v_{2}^{1,0}(B(R - \sigma_{1}R) \times [t_{0} - R^{2}, t_{0}])$$

$$\leq \gamma [(\sigma_{1}R)^{-2} + (\sigma_{2}R^{2})^{-1}] ||(v - k)^{+}||^{2} v_{2}^{2}, q_{R}^{2}$$

$$+ \gamma \left\{ \int_{t_{0}-R^{2}}^{t_{0}} [meas A_{k,R}^{+}(\tau)]^{\frac{r}{q}} d\tau \right\} = \frac{2}{r} (1 + \kappa)$$

$$+ \sup_{t \in [t_{0}-R^{2}, t_{0}]} \tilde{\zeta}_{b}^{+}(k, t_{0}-R^{2}, t)$$

for all $k \ge k_1(R)$ and satisfying (5.2).

A similar argument holds for $(x,t) \to (v-k)^{-}(x,t)$ provided that $k \le k_2(R)$. It yields inequalities to which we will refer to as $(5.1)^{-}$.

We denote with so the scallest natural number satisfying

$$\frac{2\mathbf{v}}{2\mathbf{s}} < \mathbf{s} \quad , \quad \mathbf{s} \geq 2 \quad ,$$

Let . be any positive number and construct the (BC)

$$c_{R_0}^{n-1} = x - x_0 \left\{ < R_0^{n-1} + \left(t_0 - R_0^2, t_0^{n-1}, B(R_0) + K + L \right) \right\}$$

$$c_0 = \pi R_0^2 , \quad r = \min \left\{ c_0^{n-1} + \frac{2k_1}{N+2k_1} \right\} .$$

Here $|c_0^{\dagger}|$ is the number claimed by Terma 3.1.

Lemma 5.2: Assume that

(i)
$$2M \ge -2 \text{ case } V$$
 (ii) $\frac{\lambda}{2^{s+3}} \ge \frac{N^{s}}{k_0^2}$.

 $\mathtt{The}\, n$

osc v < max{. -
$$\frac{1}{2^{s+3}}$$
; $2^{s+1} \hat{\omega}_{K}(R_{0})$ }.

Froof of Lemma 5.2: If one $v + \frac{1}{2^S}$ then the conclusion of the lemma is $v_{F_0} = v_{T_0}$ trivial. Analogously if $v + 2^{S+1} \cdot c_{K}(F_0)$, there is nothing to prove. So assume that

(5.4)
$$\frac{\operatorname{osc} v > \frac{u}{2^{\varepsilon}} > \frac{\omega}{2^{\varepsilon+1}} > \widehat{\omega}_{K}(R_{0})}{\zeta_{R_{0}}^{\omega} - T}.$$

Then at least one of the following two inequalities holds. Either

$$(5.5)_1$$
 $k_1(F_0) \cdot \frac{1}{1} - \frac{a}{2^{5+2}}$, or

$$(5.5)_2 = V_2(0_1) = V_1 + \frac{4}{2} \frac{4}{12}$$
.

Here

The second secon

$$\frac{Q_{p,q}^{*} - P_{p,q}^{*}}{Q_{p,q}^{*}} = \frac{Q_{p,q}^{*} - P_{p,q}^{*}}{Q_{p,q}^{*}} = \frac{Q_{p,q}^{*}}{Q_{p,q}^{*}} = \frac{Q_$$

Indeed if both $(5.5)_1$ and $(5.5)_2$ are violated then

$$\hat{z}_{K}^{(E_{0})} \geq k_{1}^{(F_{0})} - k_{2}^{(F_{0})} + \hat{z}^{+} - \hat{z}^{-} - \frac{1}{2^{s+1}} = \frac{1}{2^{s+1}},$$

contradicting (5.4).

Suppose that $(5.5)_2$ holds true and observe that

meas((x,t) +
$$Q_{R_0}^{(s)}$$
 | $v(x,t)$ < ... + $\frac{1}{2^{s+2}} < k_2$ \ \leq \(\frac{1}{2} \)

$$\leq \text{meas}[Q_{R_{\hat{Q}}}^{N} \cap Q_{T}^{-}] \leq c_{\hat{Q}}^{-\frac{2\kappa_{\hat{1}}}{N+2\kappa_{\hat{1}}}} \kappa_{N}^{-R_{\hat{Q}}^{N+2}} \ .$$

Consider inequalities (5.1) for $(v-k)^-$; $k \le x^- + \frac{\alpha}{2^{s+2}}$, and apply lemma

3.1. It gives the following alternative. Either

(i)
$$H = \text{ess sup } (v - (1 + \frac{N}{2} + \frac{N}{2}))^{-1} = \frac{\frac{N\kappa}{2}}{2}$$
, or

(ii) meas
$$Q_{\frac{R_0}{2}}^{-1}(\frac{1}{2} + \frac{1}{2^{s+2}} - \frac{1}{2}H) = 0$$
.

If (i) occurs then

ess inf v .
$$- + \frac{1}{2^{5+2}} = \mathbb{P}_{2}^{\frac{2}{5}}$$
.

If (ii) is valid then

ess inf
$$v > \mu^{-} + \frac{\omega}{2^{s+2}} - \frac{1}{2} \frac{\omega}{2^{s+2}}$$

$$Q_{R_{\underline{O}}}^{\omega} \cap \Omega_{T}$$

$$= \mu^{-} + \frac{\omega}{2^{s+3}}.$$

Hence in either case

$$\begin{array}{ccc}
\operatorname{osc} & v < \omega - \frac{\omega}{2^{s+3}} \\
\mathbb{Q}_{\frac{R_0}{2}}^{\omega} \cap \mathbb{T}_{T}
\end{array}$$

The lemma is proved.

<u>Lemma 5.3</u>: Let $\mathbf{x}_0 \in \text{int } K \subset \mathbb{S}$ and \mathbb{R}_0 so small that $\mathbb{B}(\mathbb{R}_0) = K$. There exists a pair of sequences $\{\mathbb{R}_n^-\} \setminus 0$, $\{\mathbb{M}_n^-\} \setminus 0$ such that

$$\begin{array}{ll} \text{osc} & v \leq M_n \\ & \\ N_n \\ Q_{R_n} & Q_{T_n} \end{array} \qquad n = 1, 2, \dots,$$

The sequences $\{R_n^-\}$ and $\{M_n^-\}$ depend only upon the data and the modulus of continuity $\hat{\mathbb{A}}_K(\cdot)$ of $x\mapsto v_0^-(x)$ in K.

Proof of lemma 5.3: For $R_1 \leq R_0$ define

$$M_1 = \max\{2M ; 2^{s+1} | \hat{\omega}_K(R_1)\} \ge 2M$$
,

and select P_1 so small that

$$\mathbb{E}_{1}^{\frac{N_{k}}{2}} = \left(\frac{2M}{2^{s+3}}\right) \simeq \left(\frac{\mathbb{E}_{1}}{2^{s+3}}\right) .$$

Then construct inductively the sequences $\{R_n^-\}$ and $\{X_n^-\}$ as follows:

$$M_{n+1} = \max\{M_n - \frac{M_n}{2^{s+3}}; 2^{s+1} \hat{\delta}_{K}(R_n)\}; n = 1, 2, ...$$

$$R_{n+1} = \min \left\{ \left(\frac{\frac{M_n}{2}}{2^{s+3}} \right)^{\frac{2}{N+1}}; \frac{R_n}{2} \right\}.$$

It is immediate to verify that $\{M_n\} \setminus 0$ and $\{R_n\} \setminus 0$. By virtue of Lemma 2.5, the conclusion of Lemma 5.3 holds true for n=1. Suppose it holds for n and let us show that it holds for n+1. By assumption osc $v \leq M_n$ and $R_n^{NK/2} \leq \frac{M_n}{2^{s+3}}$. Hence by Lemma 5.2 $Q_{R_n}^{N} \cap \Omega_T$

$$\begin{array}{c} \text{osc } v \leq \frac{M}{n+1} \\ \frac{M}{2} \\ n \\ \frac{n}{2} \end{array}.$$

To conclude the proof observe that

$$Q_{R_{n+1}}^{M_{n+1}} \cap \mathcal{A}_{T} \in Q_{R_{n}}^{M_{n}} \cap \mathcal{Q}_{T} .$$

Proof of Theorem 5.1: It is an immediate consequence of Lemma 5.3.

Corollary: Let \mathbb{C}^* be an open set contained in \mathbb{C} and assume that $\mathbf{v} \to \mathbf{v}_0^-(\mathbf{x})$ is continuous on \mathbb{C}^* with modulus of continuity $\frac{\partial}{\partial K}(\cdot)$ uniform on every compact $K \in \mathbb{C}^*$ then \mathbf{v} is continuous on $K \times [0,T]$ and

$$|v(x_1,t_1) - v(x_2,t_2)|^2 + ||x_1 - x_2|| + ||t_1 - t_2||^{\frac{1}{2}}$$

 $\forall (\mathbf{x_i}, \mathbf{t_i}) \in K = [0, T] \quad i = 1, 2 \text{ .} \quad \text{the continuous non-decreasing function} \\ + \frac{1}{K}(\cdot), \quad + \frac{1}{K}(\cdot) = 0 \quad \text{depends only upon the data and } \cdot \frac{1}{K}(\cdot) \text{ .}$

[B] The case of variational boundary data.

Consider formally the problem

(5.6)
$$\begin{cases} \frac{3}{-t} S(v) - \operatorname{div} a(x,t,v,V_{x}v) + b(x,t,v,V_{x}v) & 0 \\ a(x,t,v,V_{x}v) \cdot n_{S_{T}}(x,t) = g(x,t,v) & \text{on } S_{T} \\ v(x,0) = v_{0}(x) & v_{0}(x) \neq 0 \text{ a.e. in } \end{cases}$$

where $n_{S_T} = (n_{x_1}, n_{x_2}, \dots, n_{x_n})$ denotes the outer unit normal to S_T . We assume that a(x,t,v,p) and b(x,t,v,p) satisfy $[h_1] - [h_2]$ and that $v_0 \in L_2(\Omega)$. On the boundary data g(x,t,v) we assume that

[G] g is continuous over $S_{\frac{1}{T}} \times \mathbb{R}$ and admits an extension g(x,t,v) over $\frac{1}{T}$ such that

$$\left\| \frac{\partial}{\partial \mathbf{v}} \dot{\mathbf{g}}(\mathbf{x}, \mathbf{t}, \mathbf{v}) \right\|_{2, \mathbf{T}} < C < \infty$$

for some positive constant C.

Essentially we are imposing on $\,g\,$ a growth at most linear with respect to $\,v\,.\,$

By a weak solution of (5.6) we mean a function $v + w_2^{1,1}(c_T)$ satisfying

$$(5.7) = \int (\mathbf{x}, \tau) \cdot [\mathbf{v} \succeq 0] \mathbf{v} \, d\mathbf{x} = \begin{bmatrix} \mathbf{t} \\ \mathbf{t}_0 \end{bmatrix} \int \mathbf{v}(\mathbf{x}, \tau) \cdot [\mathbf{v} \succeq 0] = \frac{\tau}{2\mathbf{t}} \mathbf{c} \, d\mathbf{x} d\tau + \frac{\tau}{2\mathbf{t}} \mathbf{c} d\mathbf{x} d\tau + \frac{\tau}{2\mathbf{t}} \mathbf$$

$$\frac{1}{2} \left(\frac{1}{2} \left$$

 $\psi v = \psi_{q}^{(r)}(\frac{1}{r})$ and any interval $\{t_{q},t\} = [0,T]$, and $v(x_{r}) = v_{q}(x)$ in the sense of the spaces over

is continuous in (2.7) Notice for all $\mathcal{F} = \frac{\sqrt{1}}{\sqrt{2}}, \frac{1}{\sqrt{2}}$, hence a weak solution of (5.6) is also a weak solution of (1.6). Therefore by the results of scott in (2.4) we is continuous in $\frac{1}{\sqrt{2}}$. Hereever if w_i is continuous in $\frac{1}{\sqrt{2}}$ tion we is continuous allowable \mathcal{F} .

The aim of the section is to prove that the continuity of v can be extended to $v_{\overline{q}}$ and that if $v_{\overline{q}}$ is continuous in $\overline{}$, then v is continuous in $\overline{}$. In a precise way we want to prove the following theorem.

Theorem 1.7: Assume that $x = is \ a = c^2$ manifold in \mathbf{F}^{N-1} , suppose that [3] h.1.h., and let $x = w_1^{1/4} f_{-n}$) is a weak solution of (5.6) such that

For every $x \in [0, \infty]$ explicts a continuous non-lear acting function $(x \in \mathbb{R}^n) \times \mathbb{R}^n, \quad [0, \infty] \qquad \text{excitation}$

$$v(x_1,t_1) = v(x_2,t_2) \ge (x_1 - x_2 + t_1 - t_2)^{\frac{1}{2}}$$
,

Moreover if \mathbf{v}_0 is continuous over all $\overline{\mathbb{S}}$, then there exists $\mathbf{s} \to \mathbb{R}_0^+(\mathbf{s}): \mathbb{R}^+ \to \mathbb{R}^+, \ \omega_0^-(\mathbf{0}) = 0$ continuous and non decreasing such that

$$|v(x_1,t_1) - v(x_2,t_2)| \le \omega_0(|x_1 - x_2| + |t_1 - t_2|^{\frac{1}{2}})$$

for all $(x_i,t_i) \in \overline{\Omega}_T$, i = 1,2.

The functions $\omega_\sigma(\cdot)$ can be determined in dependence of the data and the positive number σ , and $\omega_O(\cdot)$ can be determined only in terms of the data and the modulus of continuity of v_O .

The proof of the theorem is essentially the same as the proof of interior regularity and is based on the same arguments of sections 3.4. The difference is that instead of working on cylinders \mathcal{Q}_{R} here we are dealing with cylindrical domains of the type

$$(B(R) \times [t_0 - R^2, t_0]$$
.

We bound ourselves to describe the differences occurring in the proof. Fix \mathbf{x}_0 , 3n and consider the portion of the boundary 3n given by

$$\mathbf{S}_0 \neq 3\Omega \; \text{n} \; \{ \left| \mathbf{x} - \mathbf{x}_0 \right| < \overline{R} \}$$
 , $\overline{R} > 0$ given .

Our arguments being local in nature, we may assume without loss of generality that S_0 lies on the hyperplane $\mathbf{x}_N=0$. Indeed this can always be achieved by a local change of coordinates in identity (5.7) written for example for test functions $\varphi(\cdot,t)$ supported in a neighborhood of \mathbf{x}_0 , for $t\in\{0,T\}$.

[B]₁. Inequalities analogous to (2.7):

Let
$$(\mathbf{x}_0, \mathbf{t}_0)$$
 , $\mathbf{S}_{\mathbf{T}}$, $\mathbf{t}_0 \ge 0$ and set

$$C_{E} = \{\{\mathbf{x} - \mathbf{x}_{0}\} \times \mathbb{R}^{2} + \mathbf{x}_{0}\}, \quad \mathbb{R} \subseteq \overline{\mathbf{k}}$$

$$C_{E} = \mathbb{R}^{2} \times [\mathbf{t}_{0} - \mathbb{R}^{2}, \mathbf{t}_{0}], \quad \mathbb{R}^{2} \times \mathbf{t}_{0}.$$

Since 3, around \mathbf{x}_0 is a portion of the hyperplane $\mathbf{x}_n = 0$ and $\mathbf{E} \subseteq \mathbb{R}$, \mathbb{E}_F is the half ball $\{ |\mathbf{x} - \mathbf{x}_0| > \mathbb{R}, \mathbf{x}_N \geq 0 \}$ and \mathbb{C}_E is the half cylinder obtained by intersecting the lateral cylinder (LC) \mathbb{Q}_E with \mathbb{T}_F . Moreover notice that since $\mathbb{R}^2 < \mathbf{t}_0$, \mathbb{C}_E does not intersect \mathbb{R}^2 , at \mathbb{T}_F at \mathbb{T}_F .

Our next task is to derive inequalities analogous to (2.7) over the domains ${\tt C}_{\tt m}$ and

$$c_{R}(\sigma_{1},\sigma_{2}) = c_{R}(\sigma_{1}R) \times [t_{0} - (1 - \sigma_{2})R^{2},t_{0}]$$
,
$$\sigma_{1},\sigma_{2} \in (0,1)$$
.

This is done by selecting in (5.7) test functions $\varphi = t(v - k)^{\frac{1}{2}} \zeta^2$, where $(x,t) \to \zeta^2(x,t)$ is chosen as in (a).

All the terms on the left hand side of (5.7) are treated as in the derivation of (2.7) except for the different domain of integration. We remark in this connection that $\zeta(x,t)$ vanishes on the parabolic boundary of Q_R , and not on the parabolic boundary of C_R .

We estimate the term involving an integration over [32] on the right hand side of (5.7) by transforming it in a volume integral as follows

$$I = \int_{t_0 - i}^{t_0} \int_{-i}^{2} g(x, i, v) (v - k)^{\frac{1}{2}} \zeta^2(x, i) dx dx = \int_{t_0 - i}^{t_0} \int_{-i}^{2} g(x, i, v) (v - k)^{\frac{1}{2}} z^2(x, i) dx dx = \int_{-i}^{i} \int_{-i}^{2} div[g(x, i, v) (v - k)^{\frac{1}{2}} z^2(x, i)] dx dx.$$

We expand the integrand, use hypothesis [G] and perform routine calculations involving the Cauchy inequality $ab \le \varepsilon^{-1}a^2 + \varepsilon b^2$, $\varepsilon > 0$, to obtain the estimate

$$I \leq \gamma_{1}(\varepsilon) \int_{C_{R}} \int (v - k)^{\frac{1}{2}} |\nabla_{\mathbf{x}} \zeta|^{2} dx d\tau +$$

$$+ \gamma_{2}(\varepsilon) \int_{C_{R}} \int \chi[(v - k)^{\frac{1}{2}} > 0] \zeta^{2}(\mathbf{x}, \tau) dx d\tau +$$

$$+ \varepsilon \int_{C_{R}} \int |\nabla_{\mathbf{x}} (v - k)^{\frac{1}{2}}|^{2} \zeta^{2}(\mathbf{x}, \tau) dx d\tau ,$$

where $\varepsilon \geq 0$ is arbitrary and γ_1 , γ_2 are constants depending upon the data and ε .

These remarks prove that there exists constants $|\gamma|$ and $|\delta|$ such that for all $|k|\in {\rm I\!R}|$ satisfying

(5.8)
$$\operatorname{ess\ sup\ }(v-k)^{\frac{+}{-}}<\delta\ ,$$

$$C_{k}$$

we have the inequalities

$$(5.9) | (\mathbf{v} - \mathbf{k})^{\pm} |_{\mathbf{v}_{2}^{1}, 0}^{2} (\mathbf{c}_{\mathbf{E}}(\mathbf{v}_{1}, \mathbf{v}_{2})) = \gamma ((\mathbf{v}_{1}^{\mathbf{E}})^{-2} + (\mathbf{v}_{2}^{\mathbf{E}^{2}})^{-1}) | (\mathbf{v} - \mathbf{k})^{\pm} |_{\mathbf{v}_{2}, \mathbf{c}_{\mathbf{R}}}^{2} +$$

$$+ + \left\{ \begin{array}{l} \frac{t}{100} \\ \frac{t}{100} \\ \frac{t}{100} + \frac{2}{100} \end{array} \right\} = \left\{ \frac{r}{100} + \frac{r}{100} + \frac{2}{100} + \frac{2}{100$$

where $\frac{1}{a}(k,t_0 = k^2,t)$ is define in $\frac{\Phi_0^{\pm}(k,t_0 + R^2,t)}{a}$ except for the different domain of integration. The parlities (5.9) hold for all k satisfying (5.8) and all $t_1,t_2 = (e,1)$.

We remark that the constants γ and δ in (5.9) might differ from the analogous constants in (2.7). This is due of course to the extra term involved in the boundary integral.

Gemma 2.1 remains unchanged and Lemma 2.2 now is stated as follows.

 $\frac{\log (n + 8.4)}{2 \log (n + 8)} = \frac{1}{2} + \frac{1}{2} + \frac{1}{2} = \frac{1}{2} + \frac{1}{2} + \frac{1}{2} = \frac{1}{2} + \frac{1}{2} + \frac{1}{2} = \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}$

$$\psi(\mathbf{x},t) = \psi(\dot{\mathbf{v}}(\mathbf{x},t)) \approx \left(n^{+} \left[\frac{n}{1 - (\mathbf{v} - \mathbf{k})^{+} + n} \right] \right)$$

then there exists a constant C depending only upon the data such that for all $t \in [t_0 + \delta \kappa^2, t_0]$

$$\int_{\mathbb{R}+\sigma_1 R} e^{2(\mathbf{x},t)d\mathbf{x}} \leq \int_{\mathbb{R}} e^{2(\mathbf{x},t)} e^{-2R^2} d\mathbf{x} +$$

$$+\frac{c}{c^{2}}(1+\ln\frac{\pi}{\eta})(1+\frac{R^{NR}}{\eta^{2}})$$
 meas C_{R} .

As remarked after Lemma 2.2, also in the present situation an analogous result holds for $(v-k)^{-}$, $k \in \emptyset$.

[B]₂. Proceeding in the proof we see that Lemma 3.1 holds in the present diffuction for the domain C_R instead of for the cylinder Q_R . The only diffication regards the proof of the recursion inequalities [I] - [II]. For those we used the embedding inequality (2.11) valid for functions of $\frac{21}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, do not vanish on the lateral boundary of $\frac{1}{2}$, therefore we must use inequality (2.12), and observe that for the domain $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, $\frac{1}{2}$, one can consider the constant in (2.12) as independent of $\frac{1}{2}$.

Finally the last modification occurs in Lemma 3.4 in the use of DeGiorgi's inequality (2.9). Now such an inequality holds also for convex domain, therefore (2.9) is valid with B(R) replaced by Ω_R . The remainder of the proof stays unchanged. The first assertion of the theorem is proved. For the second part we consider domain C_R with $t_0 - R^2 < 0$ and over them carry on the arguments of Lemma 5.2 - 5.3 with the modifications indicated above.

[C] The case of homogeneous Dirichlet boundary data.

We let $\mathbf{v} \in \mathbb{W}_2^{1,1}(\Omega_{\mathbf{T}})$ be a weak solution of (1.6) which in addition satisfies

$$v \mid S_T = 0 \quad (x,t) \in S_T$$

in the sense of the traces over S_T . In this paragraph we investigate under what assumptions on 30—the interior continuity of v—can be extended up to the lateral boundary S_T —of Ω_T . On 30—assume the following:

(P) $\exists \theta^* > 0$, $R_0 > 0$ such that $\forall x_0 \in \partial \Omega$ and every ball

B(R) centered at x_0 , $R \le R_0$,

$$meas[\Omega \cap B(R)] < (1 - e^*)meas B(R)$$
.

Theorem 5.3: Let v, $W_2^{1,1}(\mathbb{C}_T)$ be a weak solution of (1.6) such that $\|v\|_{\infty,\mathbb{C}_T} \leq \mathbb{M} < \alpha$, and $v|_{S_T} = 0$ in the sense of the traces. There exist $0 < \eta < 1$ and a constant I, such that

$$|\mathbf{v}(\mathbf{x},t)| \leq L(\operatorname{dist}[(\mathbf{x},t),3.3])^{\eta}$$
.

Moreover if $v(x,0) + v_0(x)$ in the sense of the traces over , and if $v_0 + c(\overline{L})$, $v_0|_{\partial \mathbb{D}} = 0$, then there exists a continuous non-decreasing function

 $L(\cdot)$: $\mathbb{R}^+ \cdot \mathbb{R}^+$, L(0) = 0 such that

$$|v(x_1,t_1) - v(x_2,t_2)| \le a(|x_1 - x_2| + |t_1 - t_2|^{\frac{1}{2}})$$

for all $(\mathbf{x}_i, \mathbf{t}_i) \in \overline{\mathbb{Q}}_T$ i=1,2. The numbers p and L depend uniquely upon the data, whereas $\omega(\cdot)$ can be determined in dependence of the data and the modulu, of continuity of \mathbf{v}_0 in $\overline{\mathbb{R}}$.

The theorem is a consequence of the following inequalities valid on every (LC) $|\mathcal{Q}_{\pi}\rangle$

$$(5.10) | (v - k)^{\pm} |_{V_{2}^{1,0}[C_{R}(c_{1},c_{2})]}^{2} \leq Y[(c_{1}R)^{-2} + (c_{2}R^{2})^{-1}] | | (v - k)^{\pm} |_{C_{R}}^{2} +$$

$$+\gamma \left\{ \begin{cases} t_0 \\ t_0 - R^2 \end{cases} \text{ [meas } A_{k,E}^{\frac{t}{2}}(\tau) \cap \Omega]^{\frac{r}{2}} d\tau \right\} \xrightarrow{\frac{r}{2}} (1+\kappa) \\ + \sup_{t \in [t_0 - R^2, t_0]} \left\{ t_0 - R^2, t_0 \right\}$$

for all k - E such that

(5.11) eas sup
$$(v - k)^{\frac{1}{2}} < \delta$$
.
 $C_E \cap C_T$

Here δ is the same number introduced in (2.6) and γ is the same constant appearing in (2.7). The definition of $\hat{\beta}_a(k,t_0\sim F^2,t)$ is obvious.

Inequalities (5.10) are derived from identity (2.1) upon the choice of $c = \pm (v - k)^{\pm 1/2}$, where it is selected as in (a).

Notice that since $\mathbf{v} + \hat{\mathbf{w}}_2^{1,1}(\mathbb{C}_T)$ we have also $(\mathbf{v} - \mathbf{k})^{\frac{1}{2}} + \hat{\mathbf{w}}_2^{1,1}(\mathbb{S}_T)$ and the refere such a choice of \mathcal{F} in (2.1) is justified. We will use a simplified \mathbb{R}^2 in ef.(5.15) obtained by imposing further restrictions on the levels \mathbb{R} .

If we whent $k \ge 2$ and white (5.10) for $(v-k)^+$ we see that $\frac{1}{2} h_1 d_2 = \frac{1}{2} \frac{1}{2} d_2 = 0$; poreover by looking at v = 1 extended to be zero on

that part of Q_R that remains outside Ω_T , the domains of integration in (5.10) can be replaced by $Q_R(\sigma_1,\sigma_2)$ and Q_R respectively. Hence for $(v-k)^+$, k>0 we are led to the inequalities

$$(5.10)^{+} | (v - k)^{+} |_{V_{2}^{1}, 0}^{2} (Q_{R}(\sigma_{1}, \sigma_{2})) \leq \gamma [(\sigma_{1}R)^{-2} + (\sigma_{2}R^{2})^{-1}] | | (v - k)^{+} |_{2, Q_{R}}^{2}$$

$$+ \gamma \left\{ \int_{t_{0}-R^{2}}^{t_{0}} [meas A_{k,R}^{+}(\tau)]^{\frac{r}{q}} d\tau \right\}^{\frac{2}{r} (1+\kappa)}$$

valid for all $\sigma_1, \sigma_2 \in (0,1)$, and all k > 0 satisfying (5.11).

The same argument applied to $(v - k)^{-}$, k < 0 leads to inequalities for $(v - k)^{-}$, k < 0 to which we will refer as $(5.10)^{-}$. We remark explicitly that in $(5.10)^{+}$ [(5.10 resp.] there is no restriction on the levels k other than k > 0 (k < 0 resp.) and satisfying (5.11).

Let $(\mathbf{x}_0, \mathbf{t}_0) < \mathbf{S}_T$ $\mathbf{t}_0 > 0$ be fixed and let \mathbf{R}_0 be so small that $(2\mathbf{R}_0)^2 < \mathbf{t}_0$, so that \mathbf{Q}_E , $0 \le \mathbf{R} \le 2\mathbf{R}_0$ are lateral cylinders (LC), with common "vertex" $(\mathbf{x}_0, \mathbf{t}_0)$. Set

$$\mu_{\rm cos}^{+} = \cos \sup \mathbf{v} \quad , \quad \mu_{\rm cos}^{-} = \cos \inf \mathbf{v} \quad , \quad \omega = \cos \cos \mathbf{v}$$

$$\frac{\Omega_{\rm cos}}{\Omega_{\rm cos}} = \frac{\Omega_{\rm cos}}{\Omega_{\rm$$

and, without less of generality surpose that

$$\mu^{+} > [\mu^{-}]$$
.

" recover let see I be the smallest positive integer such that

$$\frac{2N}{2^5} < \delta \quad .$$

and observe that $u^{+} - \frac{\omega}{2^{S}} > 0$. We will employ $(5.10)^{+}$ for the levels $k = u^{+} - \frac{\omega}{2^{F}}$, $p \ge s$, $p \in \mathbb{N}$, over the cylinders Q_{R} , $R \le F_{0}$.

Frame 5.5: For every $\theta_1 > 0$, there exists a positive integer p (depending on θ_1) such that either

(i)
$$\frac{\alpha}{2^p} \leq R_0^{\frac{N\kappa}{2}}$$
, or

(ii)
$$\operatorname{meas}\{(\mathbf{x},t) \in \mathcal{Q}_{R_0} \middle| \mathbf{v}(\mathbf{x},t) \geq \mu^+ - \frac{\omega}{2^p} \} \leq \theta_1 \kappa_N R_0^{N+2}$$
.

The number $\,p\,$ depends upon the data and $\,\theta_{1}^{}\,$ and it is independent of , and $\,R_{0}^{}\, \cdot\,$

Froof of Lemma 5.5: The lemma is proved in exactly the same way as Lemma 3.4. We remark that the estimate

$$\max\{B(R_0)\backslash A_{\mu}^+ + \frac{\omega}{2^s} \text{ (t)}\} > \theta^* \kappa_N R_0^N$$

for all the $[t_0 - R_0^2, t_0]$, which in Lemma 3.4 was derived from De Giorgi's inequality, in the present situation is automatic since $\frac{\partial \Omega}{\partial t}$ satisfies (P). Lemma 5.6: There exists a number $\frac{\partial \Omega}{\partial t} > 0$ such that if

$$\max\{(x,t) \in Q_{R_0} | v > \mu^+ - \frac{\omega}{2^p} \} \le \theta_1 \kappa_N R_0^{N+2}$$
,

then cities

(1) He can sup
$$(v - (v^{+} - \frac{\omega}{2^{p}}))^{+} \le \frac{\frac{N\kappa}{2}}{R_{0}^{2}}$$

or

(ii) meas{(x,t)
$$\in Q_{\frac{R_0}{2}}[v(x,t) > \mu^+ - \frac{\omega}{2^p} + \frac{1}{2}H] = 0$$
.

The number θ_1 depends only upon the data and not upon ω or R_0 .

<u>Proof of Lemma 5.6</u>: The proof is the same as for Lemma 3.1. In this case it is in fact simpler because the term $\phi_a^+(k,t_0^--R_0^2,t)=0$. It is this last fact also that makes θ_1^- independent of ω .

As a consequence of Lemmas 5.5 - 5.6 we observe the following result Lemma 5.7: Consider the decreasing sequence of numbers $\left(\frac{R_0}{2^n}\right)$ and the family of coaxial nested cylinders Q_R with common "vertex" (x_0,t_0) .

There exists a positive integer $q \in \mathbb{N}$ such that either

$$\underset{\frac{Q}{2^{n+2}}}{\text{osc}} v \leq 2^{q} \left(\frac{R_{0}}{2^{n}}\right)^{\frac{N\kappa}{2}}$$

or

$$\frac{\sum_{\substack{x \in \mathbb{N} \\ 2^{n+2}}} v \leq (1 - \frac{1}{2^{q}}) \operatorname{osc}_{\substack{x \in \mathbb{N} \\ 2^{n}}} v .$$

The first assertion of the theorem follows from Lemma 5.7. The fact that we have an estimate of Hölder type near $S_{\overline{1}}$ is a consequence of Lemma 5.7 above and Lemma 5.8 of [18] page 96-97.

The second part of the theorem is proved by estimating the escillation of v in lateral cylinders Q_R with $t_0 - R^2 \le 0$, in the same way as indicated in part [B]. We omit the details.

Femark: If $c(x,t) \equiv 0$ on $S_T^1 \subseteq S_T^-$ where S_T^1 is an open set in the relative topology of S_T^- , then the continuity can be extended up to any compact $K \subseteq S_T^1$, compact in the relative topology of S_T^- .

6. Uniform approximations:

A common device in the theory of existence of weak solutions of (1.1) subject to some initial data and to variational or Dirichlet boundary conditions, consists in solving a sequence of regularized versions of (1.1) to obtain the solution as a limit in a suitable sense of a sequence of solutions of regularized problems. It is of interest in the applications to construct the solution as a limit in the topology of the uniform convergence on compacts of $\mathbb{A}_{\mathbb{T}}$. One such application can be found in [7]. In this section we indicate how this can be realized.

Let $\mathbf{v} \in \mathbf{V}_2^{1,0}(\Omega_{\mathbf{T}})$ satisfy identity (1.5) for all $\varphi \in \mathbf{W}_2^{1,1}(\Omega_{\mathbf{T}})$ such that $\mathbf{t} + \varphi(\mathbf{x},\mathbf{t})$ has compact support in [0,T]. Suppose that there exists sequences $\{\mathbf{w}_{\mathbf{n}}\}$ and $\{\mathbf{v}_{\mathbf{n}}\} \in \{\mathbf{h}^{-1}(\mathbf{w}_{\mathbf{n}})\}$ such that

$$(6.1) \qquad \text{w}_{\text{n}}, \text{ } \text{v}_{\text{n}} \in \mathbb{W}_{2}^{1,1}(\mathbb{O}_{\text{T}})$$

$$\text{v}_{\text{n}} + \text{v} \text{ strongly in } \mathbb{L}_{2}(\mathbb{O}_{\text{T}}) \text{ and weakly in } \mathbb{V}_{2}^{1,0}(\mathbb{O}_{\text{T}})$$

$$\text{w}_{\text{n}} + \text{w} \text{ weakly in } \mathbb{V}_{2}(\mathbb{O}_{\text{T}}), \text{ w} \in \mathbb{F}(\text{v})$$

$$\text{a}_{1}(\text{x},\text{t},\text{v}_{\text{n}},\mathbb{V}_{\text{X}}\text{v}_{\text{n}}) = \mathbb{E}(\text{x},\text{t},\text{v}_{\text{n}},\mathbb{V}_{\text{X}}\text{v}_{\text{n}}) + \mathbb{E}(\text{x},\text{t},\text{v},\text{v}_{\text{X}}\text{v}) = \mathbb{E}(\text{x},\text{t},\text{v},\text{v},\mathbb{V}_{\text{X}}\text{v}) = \mathbb{E}_{2}(\mathbb{O}_{\text{T}})$$

(F.2)
$$w_n \quad \text{and} \quad v_n \quad \text{satisfy the identity}$$

$$\int_{\mathbb{R}^{N}} w_{n}(\mathbf{x}, t) \varphi(\mathbf{x}, t) \stackrel{\text{i.t.}}{\leftarrow} + \int_{\mathbf{t}_{0}}^{\mathbf{t}} \int_{\mathbb{R}^{N}} \left(-w_{n}(\mathbf{x}, t) - \frac{\partial}{\partial \mathbf{t}} \varphi(\mathbf{x}, t)\right) +$$

$$+ \left(\frac{1}{n} \cdot \mathbb{I}_{\mathbf{x}} \mathbf{w}_n + \mathbf{a}(\mathbf{x}, \mathbf{t}, \mathbf{v}_n, \mathbb{I}_{\mathbf{x}} \mathbf{v}_n)\right) + \mathbb{I}_{\mathbf{x}} \mathbf{c} + \mathbf{b}(\mathbf{x}, \mathbf{t}, \mathbf{v}_n, \mathbb{I}_{\mathbf{x}} \mathbf{v}_n) \mathbf{c} \cdot \mathbf{d} \mathbf{x} \mathbf{d} \mathbf{c} = \mathbf{0}$$

for all $\varphi \in \mathbb{W}_{2}^{1,1}(\mathbb{S}_{T})$ and all intervals $[t_{0},t] \in (0,T]$.

Since $w_n \in S(v_n)$ in the sense of the graph, (6.2) is the weak formulation of

(6.3)
$$\frac{\partial}{\partial t} \beta(v_n) - \operatorname{div} \overrightarrow{a}(x,t,v_n,\nabla_x v_n) - \frac{1}{n} \Delta \beta(v_n) + b(x,t,v_n,\nabla_x v_n) > 0 \quad \text{in } \mathcal{P}(\Omega_T) .$$

Femark: Because of the regularizing term $-n^{-1}\Delta\beta(v_n)$, the functions $(x,t)+w_n(x,t)\in\beta(v_n)$ and $(x,t)+v_n(x,t)$ are Hölder continuous over Ω_T with exponent depending upon the data and n, (see [18]).

Regularizations like (6.3) are of the type of Hopf vanishing viscosity, and were used in [2].

Furthermore we assume that the weak solution v_n of (6.3) can be obtained as a weak $w_2^{1/1}(\cdot, T)$ -limit of weak solutions of

(6.4)
$$\frac{\partial}{\partial t} \mathcal{E}_{m}(v_{n}^{m}) - \operatorname{div} \hat{a}(x, t, v_{n}^{m}, \nabla_{x} v_{n}^{m}) - \frac{1}{n} \Delta v_{n}^{m} + b(x, t, v_{n}^{m}, \nabla_{x} v_{n}^{m}) = 0 \quad \text{in} \quad \mathcal{V}(\Omega_{T})$$

where $\{f_{m}(\cdot)\}$ is a sequence of continuously differentiable regularizations of the graph $z(\cdot)$ such that

$$0 < \tau_0 \le \frac{\epsilon_m}{m}(s)$$
 $\forall s \in \mathbb{R}$

$$\theta_{m}^{*}(s) < 1 \qquad |s| \geq \frac{1}{m}$$
.

Hamely we assume that $\beta_{m}(v_{n}^{m})$, $v_{n}^{m} \in \mathbb{W}_{2}^{1,1}(\mathbb{T}_{T})$ uniformly in m and that

(i)
$$v_n^m + v_n$$
 strongly in $L_2(C_T)$, weakly in $W_2^{1,1}(C_T)$

(ii)
$$r_m(v_n^m) \to w_n + \beta(v_n)$$
 strongly in $L_2(\mathbb{T})$ and weakly in $W_2^{1,1}(\mathbb{S}_T)$

(iii)
$$a_i(x,t,v_n^m,\nabla_x v_n^m)$$
 , $b(x,t,v_n^m,\nabla_x v_n^m) \rightarrow$
$$a_i(x,t,v_n,\nabla_x v_n)$$
 , $b(x,t,v_n,\nabla_x v_n)$ weakly in $L_2(\Omega_T)$.

This second approximation is introduced only for technical reasons in order to justify the calculations below.

Theorem 6.1. Assume that ∃M < ∞ such that

$$\forall n \in \mathbb{N}$$
 $||v_n||_{\infty,\Omega_{\mathbf{T}}} \leq M$.

Then the sequence $\{v_n^{}\}$ is equicontinuous in $\Omega_T^{}$.

If $\forall n \in \mathbb{N}$ $v_n(x,0) = v_0(x) \in C(\Omega)$ in the sense of the traces over Ω , then $\{v_n\}$ is equicontinuous in $\Omega_n \cup \Omega(0)$.

If $\forall n \in \mathbb{N}$ $v_n|_{S_T} = 0$ and $v_n(x,0) = v_0(x) \in C(\overline{\Omega})$ then $\{v_n\}$ is equicontinuous in $\overline{\Omega}_m$.

Finally assume that

(i)
$$v_n(x,0) = v_0(x) \in C(\overline{\Omega})$$
 $\forall n \in \mathbb{N}$

(ii)
$$\{\vec{a}(x,t,v_n,\nabla_x v_n) - \frac{1}{n} \nabla_x v_n\} \cdot \vec{n}_{S_T} = g(x,t,v_n)$$

in the sense made precise in (5.7).

- (iii) 30 is a C^{1} manifold in \mathbb{R}^{N-1} .
- (iv) g satisfies assumptions [G] of Theorem 5.1 .

Then the sequence $\{v_n^{}\}$ is equicontinuous in $\bar{\Omega}_{\mathbf{T}}$.

<u>Proof of Theorem 6.1:</u> In section 4 we remarked that the modulus of continuity of v in \mathcal{T}_T is determined uniquely in dependence of M and the various constants appearing in (2.7) and Lemma 2.2. In view of this, to prove the

theorem will be enough to show that the functions $(v_n - k)^{\frac{1}{2}}$ satisfy inequalities like (2.7) and Lemma 2.2, with constants independent of n.

Let $(x_0,t_0) \in \Omega_T$ and R so small that $Q_R \in \Omega_T$. Let σ_1,σ_2 (0,1), construct the cylinder $Q_R(\sigma_1,\sigma_2)$ and smooth cutoff functions $(x,t) \to \zeta(x,t)$ such that

(i)
$$\zeta(x,t) \equiv 1$$
 $(x,t) \in Q_R(\sigma_1,\sigma_2)$, supp $\zeta \in O_R$,

(ii)
$$|\zeta_{t}| \leq \frac{C}{\sigma_{2}} R^{-2}$$
; $|\nabla_{\mathbf{x}} \zeta| \leq \frac{C}{\sigma_{1}} R^{-1}$

(iii)
$$\left|\Delta \zeta\right| \leq \frac{C}{\sigma_1^2} R^{-2}$$
.

It is easily checked that Lemma 2.2 carries over to the present situation with constants independent of n. We start from (6.2) choose a cutoff function $x \mapsto z(x)$ independent of t, and reproduce the same estimates in the proof of Lemma 2.2.

Mext observe that by selecting $\varphi = (v_n - k)^+ z^2$, $k \ge 0$ in (6.2) we obtain inequalities (2.7) for $(v_n - k)^+$ with $t_a^+(k, t_0 - k^2, t) = 0$. The calculations show that the constants γ and δ are independent of n (although they might differ from the analogous constants in (2.7)). The argument remains valid for $(v_n - k)^+$, $k \ge 0$. Hence we have to prove inequalities like (2.7) for $(v_n - k)^+$, $k \ge 0$ and $(v_n - k)^+$, $k \ge 0$.

For this write (6.4) in the weak form for test functions $c=(v_n^m-k)^{\frac{1}{2}-2}$. The term

$$\int_{t_0 - F^2}^{t_0} \int_{\mathbb{R}^2} \{ \hat{\mathbf{a}}(\mathbf{x}, t, \mathbf{v}_n^m, \nabla_{\mathbf{x}} \mathbf{v}_n^m) \nabla_{\mathbf{x}} \mathbf{I}(\mathbf{v}_n^m + k)^{\frac{1}{2}} \}^2 \} +$$

$$= h(\mathbf{x}, t, \mathbf{v}_n^m, \nabla_{\mathbf{x}} \mathbf{v}_n^m) (\mathbf{v}_n^m + k)^{\frac{1}{2}} \}^2 \{ \mathbf{d} \mathbf{x} \mathbf{d} \}$$

can be treated in exactly the same way as in the derivation of (2.7). Then we let $m \mapsto \infty$ (the lower semicontinuity of $\mathbb{F}_{\mathbf{X},\mathbf{n}}^{\mathbf{m}}$ in $\mathbb{F}_{\mathbf{Z}}(\mathbb{F}_{\mathbf{T}}^{\mathbf{m}})$ is employed) and observe that the constants involved are independent of n.

Next we estimate the two remaining terms

$$I_{1} = \int_{t_{0}-k^{2}}^{t_{0}} \left[-\frac{1}{2t} \beta_{m}(v_{n}^{m}) \left[-(v_{n}^{m}-k)^{\frac{1}{2}} \right] \right] L^{2}(x,t) dx dt$$

$$I_{2} = \frac{1}{n} \int_{t_{0} - \pi^{2}}^{t_{0}} \int_{\mathbb{R}^{2}} \int_{\mathbb{R}^{2}} |x_{\mathbf{x}}|^{2} \int_{\mathbb{R}^{2}} |x_{\mathbf{y}}|^{2} |\nabla_{\mathbf{x}}|^{2} (|\mathbf{v}_{\mathbf{y}}^{m}| - |\mathbf{k}|)^{\frac{\pi}{2}} |\nabla^{2}_{\mathbf{x}}(\mathbf{x}, \tau)| dx d\tau.$$

For \mathbf{I}_1 we have (we drop the subscripts \mathbf{m} and \mathbf{n} for simplicity of notation).

$$I_{1} = \int_{C_{R}} \int f'(\pm (v - k)^{\frac{1}{2}} + k) [\pm (v - k)^{\frac{1}{2}}] \frac{3}{3t} (v - k)^{\frac{1}{2}} \zeta^{2}(x, \tau) dx d\tau =$$

$$= \int_{C_{R}} \int \frac{3}{3t} \lambda [(v - k)^{\frac{1}{2}}] \cdot \zeta^{2}(x, \tau) dx d\tau ,$$

where

$$\Lambda(s) = \int_{0}^{s} \beta'(k + \xi) \xi d\xi .$$

It follows that

$$I_1 = \{ (v - k)^{\frac{1}{2}} \}^2 \xi^2(x,t) - \frac{1}{c_2 \epsilon^2} \int_{\mathbb{R}_R} \int \lambda [\pm (v - k)^{\frac{1}{2}}] dx dt .$$

The integral I_{γ} is estimated as follows:

$$nI_{2} = -\int_{Q_{R}} \int \beta(v) \nabla_{x} (v - k)^{\frac{1}{2}} \cdot \nabla_{x} \zeta^{2} dx d\tau -$$

$$-\int_{Q_{R}} \int \beta(v) (v - k)^{\frac{1}{2}} \Delta \zeta dx d\tau .$$

From this, standard calculations and limiting processes it follows that there exist constants γ , δ independent of n such that

$$(7.1) | (v_n - k)^{\frac{1}{2}} |_{V_2^{1,0}(Q_R(\sigma_1,\sigma_2))}^2 \leq \gamma [(\sigma_1 R)^{-2} + (\sigma_2 R^2)^{-1}] | | (v_n - k)^{\frac{1}{2}} |_{2,Q_R}^2 +$$

$$+ \left\{ \int_{t_0-R^2}^{t_0} \left[\operatorname{meas} A_{k,R}^{\frac{1}{2}}(\tau) \right]^{\frac{r}{q}} d\tau \right\}^{\frac{2}{r}} (1+\kappa) + \psi^{\frac{1}{r}}(k,t_0-R^2,t_0)$$

provided that ess sup $(v-k)^{\frac{1}{+}} < \delta$.

Here (\cdot,\cdot,\cdot) can be majorized by

(7.2)
$$\frac{1}{2} \leq \frac{\text{const}}{\min\{\sigma_1, \sigma_2\} R^2} \int_{\mathbb{Q}_R} \{ (v_n - k)^{\frac{1}{2}} + \chi[(v_n - k)^{\frac{1}{2}} > 0] \} dx d\tau .$$

Moreover : vanishes if (7.1) are written for $(v_n - k)^+ - k \ge 0$, or $(v_n - k)^-, k \le 0$.

The term $\frac{1}{a}$ is slightly different from the ϕ_a^{\pm} in (2.7). The only part of the proof of Theorem 1 where (2.7) has been employed with $\phi_a^{\pm} \not\equiv 0$ is Lemma 3.1. In such a lemma we estimated ϕ_a^{\pm} as

$$\frac{1}{a}(k,t_0-p^2,t) = \frac{-2v}{2^{p^2}} \int_{\mathbb{Q}_{\mathbb{R}}} \int (v-k)^{\frac{1}{2}} dx dt$$
.

By following the various steps in the proof of Lemma 3.1 it is easily checked that the extra term

$$\frac{1}{\min[\sigma_1,\sigma_2]_R^2} \int_{Q_R} \int_{\mathbb{R}} \chi[(v_n - k)^{\pm} > 0] dx d\tau$$

in (7.2) does not affect the result. A few minor changes are necessary which are left to the reader.

For the continuity up to the boundary the same arguments of section 6 are valid in the present situation. The proof is complete.

REFERENCES

- [1] Brèzis, H. On some degenerate non-linear parabolic equation, Proceedings
 AMS Vol I, XVII (1968).
- [2] Caffarelli, L. and Evans, L. C. Continuity of the temperature in the two-phase Stefan problem, (to appear).
- [3] Caffarelli, L. and Friedman, A. Continuity of the density of a gas flow in a porous medium, Trans. Amer. Math. Soc. (to appear).
- [4] Caffarelli, L. and Friedman, A. Regularity of the free boundary of a gas flow in an n-dimensional porous medium, Indiana Univ. J. of Math. Vol. 29 #3 (1980).
- [5] Cannon, J. R. and Di Benedetto, E. On the existence of solution of boundary value problems in Fast Chemical reactions, Bollettino U.M.I. (5) 15-B (1978).
- [6] Cannon, J. R. and Di Benedetto, E. On the existence of weak solutions to an n-dimensional Stefan problem with non-linear boundary conditions,

 SIAM J. on Math. Analysis Vol. 11 #4 (1980).
- [7] Cannon, J. R., Di Benedetto, E. and Knightly, G. H. The Stefan problem with convection: The non-steady state case, (in preparation).
- [8] Cannon, J. R. and Hill, C. D. On the movement of a chemical reaction interface, Indiana Math. J. 20, (1970).
- [9] Cannon, J. R. and Fasano, A. Boundary value multidimensional problems in fast chemical reactions, Archive for Fat. Mech. and Analysis Vol. 53 #1 (1973).
- [10] Cannon, J. R., Henry, D. B. and Kotlow, D. B. Classical solutions of the one-dimensional two-phase Stefan problem, Annali di Mat. Pura e Appl. (IV) Vol. CVII (1976).

- [21] De Giorgi, E. Sulla differenziabilitá e l'analiticitá delle estremali degli integrali multipli regolari, Mem. Accad. Sci. Torino. Cl. Sc. Fis. Mat. Mat. (3) 3 (1957).
- [12] Di Benedetto, E. Regularity properties of the solution of an n-dimensional two-phase Stefan problem, Bollettino U.M.I. (to appear).
- [13] Di Benedetto, F. and Showalter, R. E. Implicit degenerate evolution equations and applications, Math. Research Center Technical Summary Report.
- [14] Fasano A., and Primicerio, M. Partially saturated porous media, (to appear).
- [15] Fasano, A., Primicerio, M. and Kamin, S. Regularity of weak solutions of one-dimensional two-phase Stefan problem, Annali di Mat. Pura e Applicata (IV) Vol. CXV (1977).
- [16] Friedman, A. The Stefan problem in several space variables, Trans. Amer. Mat. Sec. 132 (1968).
- [17] Friedman, A. Partial Differential Equations of Parabolic Type, Prentice-Hall, Englewood Cliffs, N.J. 1964.
- [15] Ladyzenskaja, O. A., Solonnikov, V. A. and Ural'ceva N.N. <u>Linear and Quasi-linear Equations of Parabolic Type</u>, Amer. Math. Soc. Transl. Math. Mono, 23 Providence, FI (1968).
- [10] Dadyzenskaja, O. A. and Ural'ceva N.N. <u>Linear and Quasi-linear Elliptic</u>

 Dynations, New York Academic Press 1968.
- [20] Michs, J. L. Quelques Méthodes de Résolution des Problèmes aux Limites

 Non Lineaires, Punod Gauthier-Villars, Paris 1969.
- [21] Liens, J. L. and Jadenes F. <u>Mon-homogeneous Boundary Value Problems and</u>
 <u>Applications</u>, Vol. I Springer-Verlag, Berlin, London and New York 1972.
- 1921 Frézkov, C. N. Apriori estimates for generalized solutions of second-order elliptic and parabolic equations, Soviet Math. (1963).

- [23] Kružkov, S. N. A priori estimates of solutions of linear parabolic equations and of boundary value problems for a certain class of quasi-linear parabolic equations, Doklady Akad. Nauk 150 (1963).
- [24] Kružkov, S. N. Results concerning the nature of the continuity of solutions of parabolic equations and some of their applications. Matematicheskie Zametki, Vo. 6, #1 1969.
- [25] Kružkov, S. N. and Sukorjanskii, S. M. Boundary value problems for systems of equations of two-phase porous flow type: Statement of the problems, questions of solvability, justification of approximate methods, Math. USSR Sbornik Vol. 33 (1977) #1.
- [26] Moser, J. A new proof of De Giorgi's theorem concerning the regularity problem for elliptic differential equations, Comm. Pure Appl. Math. 13 (1960) 457-468.
- [27] Oleinik, O. A., Kalashnikov, A. S. and Yui-Lin Chzhou. The Cauchy problem and boundary problems for equations of the type of non-stationary filtration, Izvestija Akademii Nauk SSSR Ser. Mat. 22 (1958).
- [28] Rubinstein, L. The Stefan Problem, A.M.S. Translations of Mathematical Monographs Vol. 27, Providence, RI (1971).

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It is demonstrated that weak solutions of (141) in the introduction are continuous in their domain of definition. The continuity up to the boundary is also investigated.	

